#### **Twisted Bilayer Magnet Crl<sub>3</sub>**

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### Acknowledgements

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#### **Twisted Magnets Crl<sub>3</sub> Grigory Bednik Kyoung-Min Kim Do Hun Kim** (PCS-IBS) (PCS-IBS) (KAIST) (KAIST)



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#### **3D Twisted Superconductivity** & Quasicrystal

**SungBin Lee** (KAIST)

Yong Baek Kim (Toronto)



Ref: K.-M. Kim, D. H. Kiem, G. Bednik, M. J. Han, MJP, arXiv:2206.05264 (2022)

#### **Twisted Bilayer Graphene**







**ChangHwan Yi** 

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(PCS-IBS)

**Hee Chul Park** (PCS-IBS)



#### Hofstadter **Moire Replica HOTI**

Sun-Woo Kim (KAIST/SKKU BRL→Cambridge)



## **Physics of Length Scale**

#### <u>Solid state</u> <u>lattice</u>

Magnetic Domains

Complex Network



## Length Scale in solid state





In this talk, we generalize moire materials to spin systems, "twisted bilayer magnetism"

## **Experimental progress**

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#### oLocal (AFM) measurement

Letter Published: 29 November 2021

#### Coexisting ferromagnetic–antiferromagnetic state in twisted bilayer $Crl_3$

Yang Xu, Ariana Ray, Yu-Tsun Shao, Shengwei Jiang, Kihong Lee, Daniel Weber, Joshua E. Goldberger, Kenji Watanabe, Takashi Taniguchi, David A. Muller, Kin Fai Mak 🖂 & Jie Shan 🖂

Nature Nanotechnology 17, 143–147 (2022) Cite this article

6584 Accesses | 1 Citations | 12 Altmetric | Metrics

REPORT MAGNETISM

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#### Direct visualization of magnetic domains and moiré magnetism in twisted 2D magnets

TIANCHENG SONG (D), QI-CHAO SUN (D), ERIC ANDERSON (D), CHONG WANG, JIMIN QIAN, TAKASHI TANIGUCHI (D), KENJI WATANABE (D), MICHAEL A. MCGUIRE (D), RAINER STÖHR (D), [...] XIAODONG XU (D) +4 authors Authors Info & Affiliations

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# Twisted bilayer magnet

### **Transition metal trihalides**



Honeycomb magnet
<u>Crl3</u>







 $H = \sum_{\langle i,j \rangle} J \mathbf{S}_i \cdot \mathbf{S}_j + \sum_{z_j = z_i + d} J_{ij}^{\perp} \mathbf{S}_i \cdot \mathbf{S}_j + D_z \sum_i (S_i^z)^2$ 

Sivadas et al. Nano Letters (2018)

# **Symmetry of Twisted Crl<sub>3</sub>**



#### ➤ Graphene





≻ Crl<sub>3</sub>

- Monolayer preserves  $\mathbf{C}_{\mathbf{2z}}$  and  $\mathbf{P}$  symmetry
- $C_{2z}$  is preserved in twisted bilayer
- Point group D<sub>6</sub>

- Non-magnetic I atoms break  $C_{2z}$
- Twisted bilayer breaks both  $\mathbf{C}_{\mathbf{2z}}$  and  $\mathbf{P}$  symmetry
- Point group D<sub>3</sub>

## **Ab-initio model construction**



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### Local stacking structure



## Local stacking structure



- Local FM and AFM interlayer coupling coexists.
- <u>AB sublattice symmetry breaking.</u>

### **Monte Carlo Simulations**



# **Skyrmion without DMI**

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#### Skyrmions in moire superlattice



#### Standard recipe for skyrmion :

Exchange + DMI + Magnetic field

#### In twisted bilayer magnets :

Exchange + Modulating interlayer coupling(sublattice breaking)

$$N = \int \vec{n} \cdot \left(\frac{\partial \vec{n}}{\partial x} \times \frac{\partial \vec{n}}{\partial y}\right) dA$$

# **Small-angle limit**







<u>Central observation of moire pattern</u> "as  $\theta \rightarrow 0$ , moire size diverges"

(Interlayer coupling) X (Area) VS (Intralayer coupling) X (Length) ~Collinear order



# **Magnetic phase transition**



#### Landau theoretical description

Free energy functional:

$$\begin{split} F[\mathbf{n}_t, \mathbf{n}_b] &= \sum_{l=t, b} \int d^2 \mathbf{x} \left\{ \frac{3a_0^2}{2} J[\nabla_{\mathbf{x}} \mathbf{n}_l(\mathbf{x})]^2 - D_z [n_l^z(\mathbf{n}_l)]^2 \right\} \\ &+ \int d^2 \mathbf{x} \ \bar{J}_{\perp} \ \mathbf{n}_t(\mathbf{x}) \cdot \mathbf{n}_b(\mathbf{x}), \end{split}$$

Continuum Ansatz:

$$\mathbf{n}_t = (\sin \Phi_t, 0, \cos \Phi_t), \\ \mathbf{n}_b = (-\sin \Phi_t, 0, \cos \Phi_t),$$

Expansion:

$$F[\Phi_0] = N_{\text{ncd}}(\theta)(\bar{J}_\perp - 2D_z) + \frac{a}{2}[J - J_c(\theta)]\Phi_0^2 + \frac{b}{4}J_c(\theta)\Phi_0^4 + \mathcal{O}(\Phi_0^6) \Phi_0 = \pm \sqrt{(a/b)[1 - J/J_c(\theta)]}$$



Conventional second order phase transitions as a function of tilt angle

## Magnetic phase transition II



#### Landau theoretical description

Continuum Ansatz:

$$\mathbf{n}_t = (\sin \Phi_t, 0, \cos \Phi_t), \qquad \mathbf{n}_t = (\sin \Phi_t, 0, \cos \Phi_t), \mathbf{n}_b = (-\sin \Phi_t, 0, \cos \Phi_t), \qquad \mathbf{n}_b = (0, 0, 0),$$

Expansion:

$$F[\Phi_{0}] = N_{\text{ncd}}(\theta)(\bar{J}_{\perp} - 2D_{z}) + \frac{a}{2}[J - J_{c}(\theta)]\Phi_{0}^{2} + \frac{b}{4}J_{c}(\theta)\Phi_{0}^{4} + \mathcal{O}(\Phi_{0}^{6}) F[l] = \frac{aJ\pi^{2}}{4}\left(\frac{2R}{l} - 1\right) - D_{z}N_{md}(\theta)\left[1 + \frac{1}{2}\left(1 - \frac{l}{R}\right)^{2}\right] - \bar{J}_{\perp}N_{md}(\theta)\left[\left(1 - \frac{l}{R}\right)^{2} - \frac{4}{\pi^{2}}\left(\frac{l}{R}\right)^{2}\right], \quad (4.6)$$



## **Competing scales of moire magnet**







#### Magnetic phases

#### Phase transitions

• Excitations

### **Holestein-Primakoff Boson**



Global magnetic ground state



Local harmonic oscillator



## **Dirac magnons**

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Dirac magnons are protected by coexistence of the following three symmetries.

- <u>U(1)c symmetry</u>
   (Collinearity)
- <u>U(1)v symmetry</u>
   (Valley decoupling)

<u>C2z symmetry</u>
 (Lateral shift)

Park et al. PRB (2021)

# **Topological magnons**

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MJP, Youngkuk Kim, Gil Young Cho, SungBin Lee, PRL (2019)

## Magnon phase diagram

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#### Different magnon gaps are realized as a function of twist angles.



#### **Chalker-Coddington network**



## **Overall Structure of moire magnets**

#### Magnetic phases





#### **Future Research Directions**

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#### We extend theory of moire magnetism to various magnetic materials

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<u>Development of</u> <u>Extensive Monte-Carlo methods</u>



Various magnetic materials : Spin liquid  $\alpha$ -RuCl<sub>3</sub> Magnetic TMDCs



arXiv:2206.05264 (2022)

#### **Twisted trilayer Magnet**

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Stacking dependent couplings



(In preparation)

### **Twisted triple layer**



Controlling geometry:



(In preparation)

## **Twisted triple layer**



C3 symmetry breaking order:









Analogy with Josephson junction network:







### Summary

