Induced Superconductivity in the Fractional Quantum Hall Edge in Graphene Heterostructures



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Inducing Superconducting Correlation in Quantum Hall Edge States G.-H. Lee, K.-F. Huang, D. K. Efetov, D. S. Wei, S. Hart, T. Taniguchi, K. Watanabe, A. Yacoby, P. Kim Nature Physics 13, 693–698 (2017)

Imaging Andreev Reflection in Graphene S. Bhandari, G.-H. Lee, K. Watanabe, T. Taniguchi, P. Kim & R. M. Westervelt Semiconductor Science and Technology 35, 09LT02 (2020)

Induced superconductivity in the fractional quantum Hall edge Ö. Gül, Y. Ronen, S. Y. Lee, H. Shapourian, J. Zauberman, Y. H. Lee, K. Watanabe, T. Taniguchi, A. Vishwanath, A. Yacoby, P. Kim Phys. Rev. X 12 021057 (2022)

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Exchange Statistics of Quasi-particles

Indistinguishable particles

3-Dimension (and higher)

bosons
$$\psi_{(r_1,r_2)} \rightarrow + \psi_{(r_2,r_1)}$$

fermions
$$\psi_{(r_1,r_2)} \rightarrow -\psi_{(r_2,r_1)}$$

2-Dimension is special!

For non-degenerate ground state



 $\theta = \pm \pi/m$

 $m = 0, 1, 2, 3 \dots$

If the ground state is degenerated,

$$\vec{\psi}_{(r_1,r_2)} \to U_{12}\vec{\psi}_{(r_2,r_1)}$$

 U_{12} : unitary operator

Non-Abelian Anions For Topologically Protected Qubit

Non-abelian anyons

$$\vec{\psi}_{(r_1,r_2)} \rightarrow U_{12}\vec{\psi}_{(r_2,r_1)}$$

$$U_{12} \ U_{21} \neq U_{21} \ U_{12}$$

By braiding the anyons one can create non-local entangled qubits



Das Sarma, Freedman, Nayak Physics Today (2006)

Topologically protected quantum computing with non-abelian anyons



Field and Simula, (2018)

Quantum Material Platforms For Topologically Protected Qubit

5/2 Fractional Quantum Hall State



theory:

Moore, G. & Read, N., Nucl. Phys. B 360, 362 (1991) Das Sarma, S. et al. PRL 94, 166802 (2005) Stern, A. & Halperine, B.I., PRL 96, 016802 (2006) Nayak, C. et al. Rev. Mod. Phys. 80, 1083 (2008) experiment

Dolev, M. et al. Nature 452, 829 (2008) Bid, A. et al. Nature 466, 585 (2010) Mitali B. et al. Nature 559, 205 (2018) Willett et al PNAS 2009 (Bell Labs)

Non-abelian quasiparticles in **topological superconductors** theory: Fu, Kane PRL 2008

experiment:

Hart et al Nature Phys 2014 (Harvard/Würzburg) Wiedenmann et al Nature Comm 2016 (Würzburg) Sun et al PRL 2016 (Shanghai) Li et al Nature Mat 2018 (Twente)

theory: Lutchyn et al PRL 2010; Oreg et al PRL 2010 experiment:

Mourik et al Science 2012 (Delft) Rokhinson et al Nature Phys 2012 (Purdue) Das et al Nature Phys 2012 (Weizmann) Deng et al Nano Lett 2012 (Lund/Beijing) Albrecht et al Nature 2016 (Copenhagen)

theory: Nadj Perge et al PRB 2013 experiment:

Nadj Perge et al Science 2014 (Princeton) Pawlak et al npj Quantum Info 2016 (Basel) Ruby et al PRL 2015 (Berlin)





+



1D semiconductor

topological insulator

atomic chain

Superconductor

Rise of 2D van der Waals Systems



- Semiconducting materials: WSe₂, MoSe₂, MoS₂, ...
- Complex-metallic compounds : TaSe₂, TaS₂, ...
- Magnetic materials: Fe-TaS₂, CrI₃, CrGeTe₃,...
- Superconducting: NbSe₂, WTe₂, Bi₂Sr₂CaCu₂O_{8-x}, ZrNCl,...



Pick-up Technique and Edge Contacts for Multilayer vdW Stacking



L. Wang et al, Science (2013)

Graphene Based Quantum Electronic Devices



Ronen*, Werkmeister* et al., Nature Nano (2021)

Induced Superconductivity in Fractional Quantum Hall Edge



Superconductivity

Fractional quantum Hall



Superconducting Proximity in Quantum Hall State

PHYSICAL REVIEW B 82, 184516 (2010)

Chiral topological superconductor from the quantum Hall state

Xiao-Liang Qi,^{1,2} Taylor L. Hughes,^{1,3} and Shou-Cheng Zhang¹



Proximized quantum (anomalous) Hall states:

Localized majorana zero mode in the vortex core, and delocalized majorana edge states.



Parafermions for Universal Topological Quantum Computing

 $|n\rangle$:

Localizing parafermions

PHYSICAL REVIEW X 2, 041002 (2012)

Fractionalizing Majorana Fermions: Non-Abelian Statistics on the Edges of Abelian Quantum Hall States

Netanel H. Lindner Institute of Quantum Information and Matter, California Institute of Technology, Pasadena, California 91125, USA Department of Physics, California Institute of Technology, Pasadena, California 91125, USA

> Erez Berg Department of Physics, Harvard University, Cambridge, Massachusetts 02138, USA

Gil Refael Department of Physics, California Institute of Technology, Pasadena, California 91125, USA

Ady Stern Department of Condensed Matter Physics, Weizmann Institute of Science, Rehovot 76100, Israel (Received 1 May 2012; published 11 October 2012) PHYSICAL REVIEW B 86, 195126 (2012)

Superconducting proximity effect on the edge of fractional topological insulators

Meng Cheng

PHYSICAL REVIEW B 87, 035132 (2013)

Fractional topological superconductor with fractionalized Majorana fermions

Abolhassan Vaezi*

 $\sigma \times \sigma = 1 + V_1 + V_2 + \dots + V_{2m-1}$

PHYSICAL REVIEW X 4, 011036 (2014)

Universal Topological Quantum Computation from a Superconductor-Abelian Quantum Hall Heterostructure

Roger S. K. Mong,¹ David J. Clarke,¹ Jason Alicea,¹ Netanel H. Lindner,¹² Paul Fendley,³ Chetan Nayak,^{3,4} Yuval Oreg,⁶ Ady Stern,⁶ Erez Berg,⁶ Kirill Shtengel,⁷⁸ and Matthew P. A. Fisher³ Braiding parafermions for quantum gating

 $=\frac{\pi}{2m}$

 s_n





v = 1/m g < 0

Proximitizing Quantum Hall Edge by Superconductors



Andreev Reflection

A.F. Andreev, Soviet Physics JETP 19, 1228 (1964)



G.E. Blonder, M. Tinkham, and T. M. Klapwijk (1982)



Solving Bogoliubov-de Gennes Equation

Crossed Andreev Reflection in Quantum Hall Edge



Goal : Measuring negative potential on the downstream edge state

Crossed Andreev Reflection in Graphene Integer Quantum Hall States



Width Dependence of CAR: Superconducting Coherence Length



G.-H. Lee *et al*. Nature Physics (2017)

600

Toward Fractional Quantum Hall CAR

Graphene Device with superconductor & graphite gates



narrow NbN superconductor



~100 nm

etched under the superconductor

graphene boron nitride graphite SiO₂ substrate

normal leads



Cross-Andreev Reflection Probability

Simultaneous Resistance measurement:



In the quantum Hall regimeLongitudinal
magneto
$$R_{xx} = V_{xx} / I \longrightarrow 0$$
Hall $R_{xy} = V_{xy} / I \longrightarrow h/e^2 v$ Cross-
Andreev $R_{CAR} = V_{CAR} / I \longrightarrow -R_{xy} / 2 < R_{CAR} < 0$

Probability of CAR: $p_{CAR} = -R_{CAR}/(R_{CAR} + R_{xy}) \sim -R_{CAR}/R_{xy}$

Gül, Ronen, et al., arXiv:2009.07836

Crossed Andreev Reflection Probability



Gül, Ronen, et al., arXiv:2009.07836

p-wave coupled Cross Andreev Edge States





v=2 edge states are just two copies of v=1.

Strong Development of CAR in Fractional Quantum Hall States



Fractional Crossed Andreev Reflection Probability



Electron Cooper Pairs versus Quasi Particle Andreev Pairs



FQH edge–Superconductor Interface: Microscopic Picture



What are the next steps?

Parafermion Josephson Junction





Tunneling into the zero mode



Quantum Hall edge to parafermion zero mode

J. Shi et al. unpublished



Outlook





Increasing CAR probability

- *atomically close distance* (~1 nm) *between counterpropagating edges*
- pairing does not rely on spin-orbit coupling
- may allow supercurrent carried by fractional charges

Electric field-control of fractional ground states

Ideas from Onder Gul et al

Summary

- CAR of IQH/FQH is observed in graphene
- Neither valley nor spin are conserved in CAR for IQHE with CAR probability is estimated to be $\sim 3 \%$
- Strong CAR for some FQH edge with up to 10% CAR probability









Sean Hart







T. Taniguchi, K. Watanabe













Ashvin Vishwanath



Crossed Andreev Reflection for IQH/FQH: Bias Dependence



Crossed Andreev Reflection for IQH at Low Magnetic Fields



Crossed Andreev Reflection for IQH/FQH: Temperature Dependence

