

Tkachenko wave: From modern field theory viewpoint

Dung Xuan Nguyen¹, Yi-Hsien Du², Sergej Moroz³ and Dam Thanh Son²

1 Institute for Basic Science (IBS), Center for Theoretical Physics of Complex Systems (PCS), Daejeon 34126, Republic of Korea.

2 Kadanoff Center for Theoretical Physics, University of Chicago, Chicago, IL 60637, USA.

3 Department of Engineering and Physics, Karlstad University, Karlstad, Sweden.



Motivation

- Higher-rank symmetries* play an important role in high-energy and condensed matter physics. One example, the *dipole symmetry*, induces UV/IR mixing and restricts the mobility of fractonic excitations.
- Vortex lattice in a rotating superfluid is a quantum crystal that exhibits a *noncommutative* version of *dipole symmetry*.
- The collective oscillations of the vortex lattice, known as the Tkachenko mode, play the role of a Goldstone boson shared by the spontaneously broken dipole and charge symmetries.
- We developed a noncommutative effective field theory of the Tkachenko mode that predicts a *cubic dependence* of the decay width of the Tkachenko mode on its energy.

Lifshitz model

- Effective theory of vortex lattice in dual description

$$\mathcal{L} = \frac{n_0}{2} n_s \epsilon^{ij} u_i \dot{u}_j + \dot{\mathbf{u}}^2 - G(\partial_i u_j + \partial_j u_i)^2 + \mathbf{u} \cdot \mathbf{e} + \frac{1}{2} (\mathbf{e}^2 + \mathbf{b}^2)$$

- n_0 and n_s are vortex density and superfluid density respectively, G determines the elastic energy.

- Dual electric field and dual magnetic field related to superfluid current and density respectively $e_i = \partial_t a_i - \partial_i a_0 = \epsilon^{ij} j_s^j$, $b = \epsilon^{ij} \partial_i a_j = n_s$, vortices are charged under the emergent $u(1)$ gauge.

- At the low energy limit, $\omega \ll 2|\Omega|$, the power counting determined by the dispersion relation $\omega \sim k^2$, $\mathbf{e}^2$ and $\dot{\mathbf{u}}^2$ terms are ignored.

- $\frac{\delta S}{\delta a_0} = \nabla \cdot \mathbf{u} = 0 \rightarrow u^i = \epsilon^{ij} \partial_j \phi$, the special phonon only has the transverse polarisation.

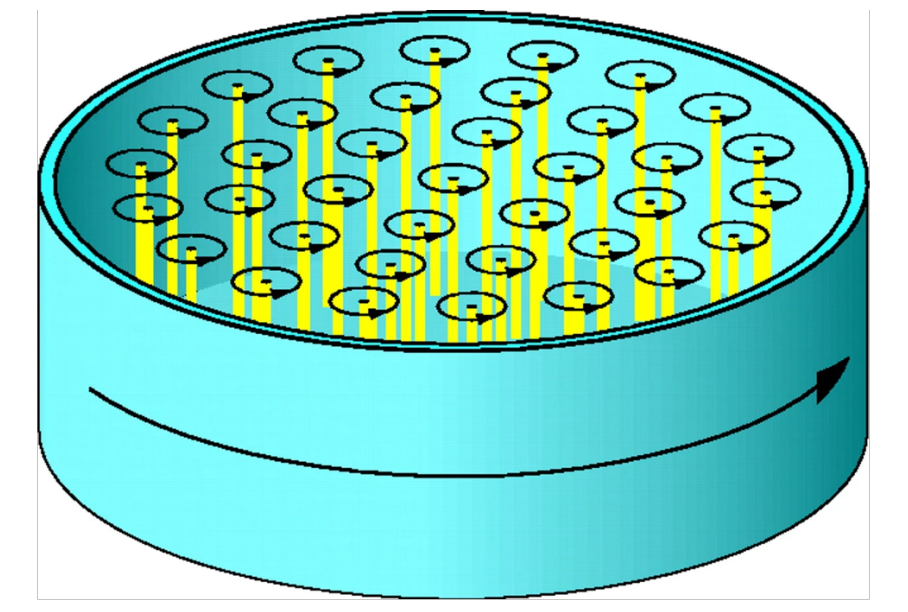
- Integrating out the dual gauge field a_μ gives the Lifshitz action

$$\mathcal{L} = \dot{\phi}^2 - c^2 \left(\nabla^2 \phi \right)^2$$

- The Lifshitz model describes the system with dipole global symmetry $\phi \rightarrow \phi + \vec{\alpha} \cdot \vec{x}$.

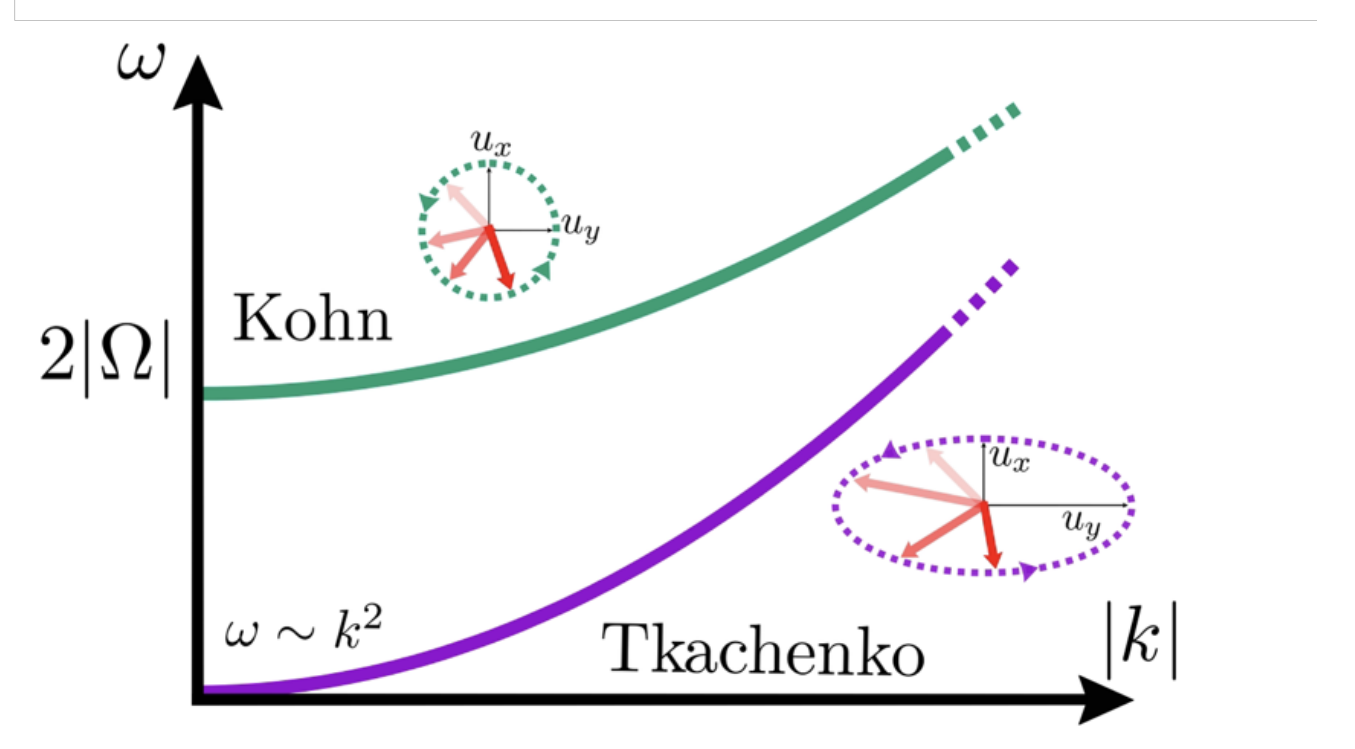
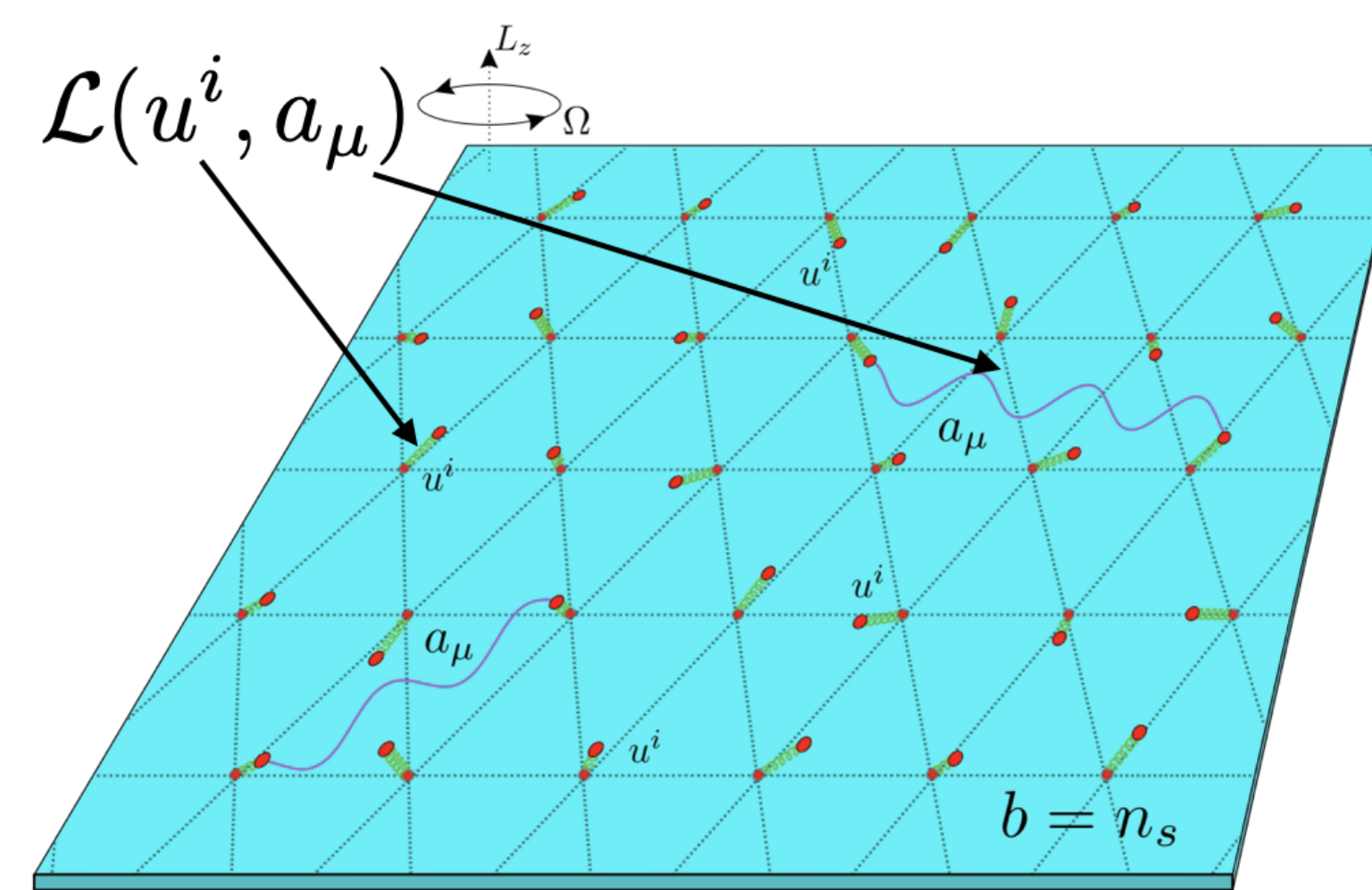
The Tkachenko wave

- Quantum vortices in superfluids have quantized circulation in units of the Planck constant.
- Vortex lattice is formed in a rotating superfluid where the vortices are distributed in a triangular lattice.
- Tkachenko way is the vibration of the vortex lattice with strange behavior:
 - Quadratic dispersion $\omega = ck^2$.
 - Only one polarization: transverse.
 - One Nambu-Goldstone boson for breaking 2-translation symmetries and $U(1)$



Circulation around a vortex line

$$\kappa = \oint d\mathbf{l} \cdot \mathbf{v}_s = \frac{\hbar}{m_{He}} \oint d\mathbf{l} \cdot \nabla \phi = \frac{h}{m_{He}}$$



Dynamical degree of freedom of the vortex lattice and Tkachenko mode.

- u^i denotes the displacement of vortices, a_μ denotes the dynamical degree of the superfluid in the *dual description* with the replacement $j_s^\mu = \epsilon^{\mu\nu\lambda} \partial_\nu a_\lambda$.
- Ω is the angular velocity of the rotating container.
- With the low energy less than the energy gap of Kohn mode, $\omega \ll 2|\Omega|$, Tkachenko mode is the only dynamical degree of freedom.
- Tkachenko wave is the **combination** of the lattice vibration (**phonon**) and the **dual photon**.

Non-commutative field theory of Tkachenko wave

- Deformed lattice described by Lagrangian coordinate

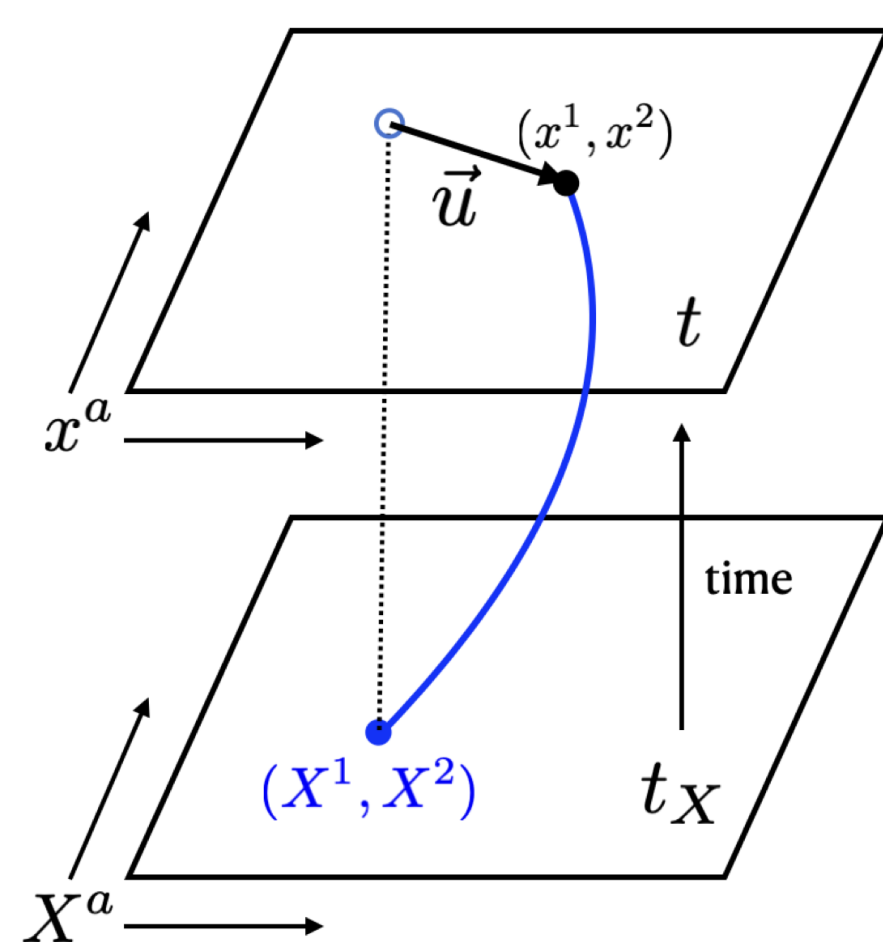
$$x^i \leftrightarrow X^a = \delta_i^a x^i - u^a$$

- Vortex current in the Lagrangian coordinate formalism

$$j_b^\mu = \frac{n_0}{2} \epsilon^{\mu\nu\lambda} \epsilon^{ab} \partial_\nu X^a \partial_\lambda X^b$$

- The low energy dynamics is constrained by the area-preserving mapping (incompressibility condition)

$$\frac{\delta S}{\delta a_0} = 0 \rightarrow \frac{1}{2} \epsilon^{ij} \epsilon^{ab} \partial_i X^a \partial_j X^b = 1$$



Non-commutative coordinate

- The angular velocity breaks time reversal T and plays the role of the magnetic field in the rotating frame.
- In the lowest Landau level projected Hilbert space, $\omega \ll 2|\Omega|$, the geometry becomes non-commutative due to the *Aharonov-Bohm phase*

$$[\hat{x}, \hat{y}] = -i\ell^2 = i\theta, \quad \ell^{-2} = B = 2\Omega m_{He}$$

- The *non-commutative* version of the *area-preserving mapping*

$$[\hat{X}, \hat{Y}] = i\theta$$

- The area-preserving can be satisfied by an unitary transformation using a non-commutative scalar field $\hat{\phi}$ as the generator

$$\hat{X}^a = e^{i\hat{\phi}} \hat{x}^a e^{-i\hat{\phi}}$$

- Tkachenko mode = non-commutative scalar field $\hat{\phi}(\hat{x})$.**
- Under the global magnetic translation ($\hat{\phi}(\hat{x})$ transformed as a dipole symmetry)

$$e^{i\vec{P} \cdot \vec{c}} \hat{\phi} \rightarrow \hat{\phi} + \vec{\alpha} \cdot \hat{x} + \frac{1}{2} \vec{c} \cdot \vec{\nabla} \phi \quad (c^i = \ell^2 \epsilon^{ij} \alpha_j)$$

- Under the $U(1)$ transformation, $\hat{\phi}$ transformed as superfluid phase $\hat{\phi} \rightarrow \hat{\phi} + \lambda$
- $\hat{\phi}$ is the shared Goldstone boson of magnetic translations and $U(1)$ global symmetries**

Summary and conclusion

- We developed a non-commutative field theory model for vortex lattice in the lowest Landau level limit.
- Nambu Goldstone bosons counting can be more complicated than the number of broken symmetry generators.
- Tkachenko is a NGB of a non commutative field theory.
- Tkachenko is a stable quasi particle $\tau \sim \omega^{-3}$.

Moyal product

- Non-commutative field theory can be mapped to a familiar commutative field theory by replacing usual product by **Moyal product**

$$\hat{A} \hat{B} \rightarrow A \star B = A \exp\left(\frac{i}{2} \epsilon^{ij} \overleftarrow{\partial}_i \overrightarrow{\partial}_j\right) B = AB + \frac{i}{2} \theta \epsilon^{ij} \partial_i A \partial_j B + \dots$$

- We trade non-commutativity for non-locality.

- Physics motivation:** Since the **coordinates do not commute**, one can't define the position of operator, **operators are blur objects**.

Non-linear effective theory and Tkachenko decay rate

- Non-commutative field theory

$$X^a = e^{i\phi} \star x^a \star e^{-i\phi} = x^a + \theta D_i \phi, \quad D_\mu \phi = -i \left(\partial_\mu e^{i\phi} \right) \star e^{-i\phi}$$

- The symmetry under magnetic translations constraints the action $\mathcal{L} = \mathcal{L}(D_t \phi, \partial_i D_j \phi)$
- Non-linear action up to cubic order of ϕ with also the rotational symmetry

$$S = \int dt d^2 \mathbf{x} \left[\dot{\phi}^2 - c^2 \left(\nabla^2 \phi \right)^2 + g_1 \dot{\phi}^3 + g_2 \dot{\phi} \left(\nabla^2 \phi \right)^2 + g_3 \left(\nabla^2 \phi \right)^3 \right]$$

- Simple dimensional analysis

$$[t] = [x^2] = -2, \quad [\omega] = [k^2] = 2 \rightarrow [g_i] = -2$$

Fixes the Beliaev decay of Tkachenko quanta

$$\Gamma \sim g^2 \omega^3$$

$$[\Gamma] = [1/t] = 2$$

- Tkachenko is a well defined excitation at low energy**

References :

- V. K. Tkachenko, Elasticity of vortex lattices, Sov. Phys. JETP **29**, 945 (1969).
- E. B. Sonin, JETP Lett. **98**, 758 (2014).
- D. X. Nguyen, A. Gromov, and S. Moroz, SciPost Phys. **9**, 076 (2020).
- Y.-H. Du, U. Mehta, D. X. Nguyen, and D. T. Son, SciPost Phys. **12**, 050 (2022).
- Yi-Hsien Du, Sergej Moroz, Dung Xuan Nguyen, Dam Thanh Son, arXiv:2212.08671 (2022).