Free energy of broadened Landau levels of a two-dimensional electron gas

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ABSTRACT

where

and

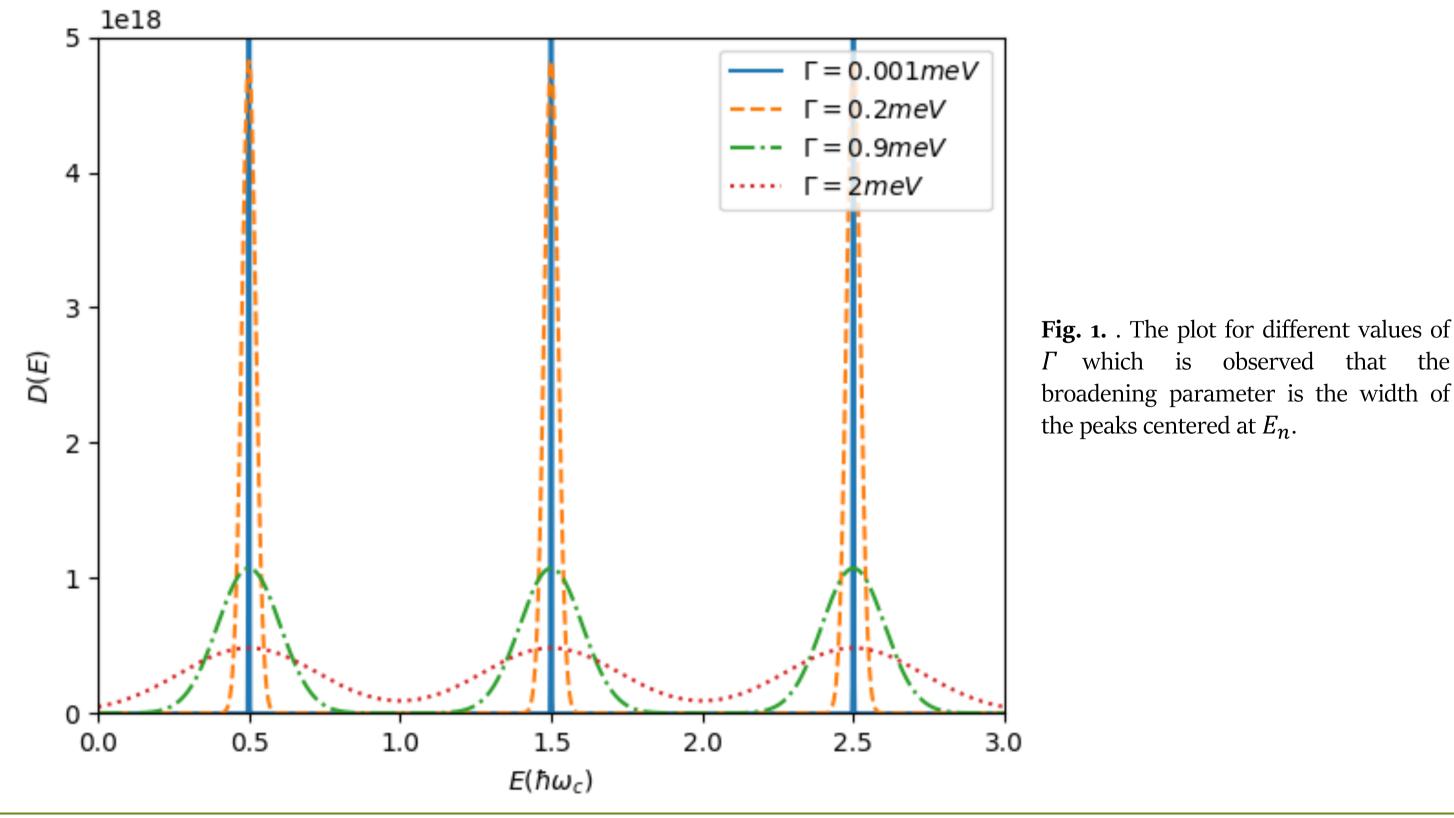
- A broadened Landau level is considered.
- An analytic expression is derived for the free energy using a Gaussian density of states (DOS).

INTRODUCTION

- Two-dimensional electron gas (2DEG) is a system where the motion of electrons is confined in two directions. [1]
- To evaluate the remaining integral, the Fermi-Dirac distribution function can ullet $d \sim d h$ be expa
- The density of states (DOS) for a broadened 2DEG under a perpendicular magnetic field which is given as

$$D(E,B) = \frac{2eB}{h} \sum_{m} \left(\frac{1}{2\pi\Gamma^2}\right)^{1/2} \exp\left[-\frac{(E-E_n)^2}{2\Gamma^2}\right]$$
(1)

• Here, $E_n = \left(n + \frac{1}{2}\right) \hbar \omega_c$ is the nth Landau level, $\omega_c = \frac{eB}{m^*}$ is the cyclotron frequency for a given effective mass m^* , and Γ is the broadening parameter.



be expanded by geometric series [2]

$$f(E) = \begin{cases} \sum_{k=0}^{\infty} (-1)^k \exp[\beta k(E-\mu)] & E < \mu \\ \sum_{k=0}^{\infty} (-1)^k \exp[\beta (k+1)(\mu-E)] & E > \mu \end{cases}$$
(6)
where the remaining integral in eq. (5)

$$J_n = \frac{\beta}{2} \int_0^{\mu} 2\nu_n dE + \frac{\beta}{2} \sum_{k=1}^{\infty} (-1)^k (\mathcal{I}_{nk} - \mathcal{J}_{nk}) \qquad (7)$$
where

$$\mathcal{I}_{nk} = \int_0^{\mu} 2\nu_n \exp[\beta k(E-\mu)] dE \qquad (8)$$
and

$$\mathcal{J}_{nk} = \int_{\mu}^{\infty} 2\nu_n \exp[\beta k(\mu - E)] dE$$
(9)

- The first term of eq. (7) can be evaluated by using integral no. 5.41 of Ref. \bullet [3]. Eqs. (8) and (9) can be evaluated by IBP simultaneously, and using the integral no. 1, 2.33 of Ref. [3].
- Upon evaluating all of the integration limits for eq. (8) and eq. (9),

$$\mathcal{I}_{nk} - \mathcal{J}_{nk} = \frac{e^{-\beta k\mu}}{\beta k} \left\{ \operatorname{erf}\left(\sqrt{\frac{1}{2\Gamma^2}} E_n\right) - \exp\left(\frac{(\Gamma\beta k)^2}{2}\right) \times (10) \right\}$$

FREE ENERGY

The free energy has the form, \bullet

$$F = \mu N - \frac{1}{\beta} \int_0^\infty D(E, B) \ln\{1 + \exp[\beta(\mu - E)\} dE$$
(2)
where $\beta = \frac{1}{k_B T}$.

• Substitute D(E) of eq. (2) by eq. (1). The compressed form of the free energy,

$$F = \mu N - \frac{2eB}{\beta h} \sum_{n} I_n \tag{3}$$

where

$$I_n = \int_0^\infty \left(\frac{1}{2\pi\Gamma^2}\right)^{1/2} \exp\left[-\frac{(E-E_n)^2}{2\Gamma^2}\right] \ln\{1 + \exp[\beta(\mu - E)]\} dE \quad (4)$$

The integral I_n can be evaluated by integration by parts (IBP) which leads to

$$I_n = uv_n \Big|_0^\infty + \frac{\beta}{2} \int_0^\infty 2v_n f(E) dE$$
(5)
where $u = \ln\{1 + \exp[\beta(\mu - E)], v_n = \frac{1}{2} \operatorname{erf}\left(\sqrt{\left(\frac{1}{2\Gamma^2}\right)} (E - E_n)\right)$ and

$$\left[\operatorname{erf}\left(\sqrt{\frac{1}{2\Gamma^2}} \left(\mu - E_n - \Gamma^2 \beta k \right) \right) + \operatorname{erf}\left(\sqrt{\frac{1}{2\Gamma^2}} \left(E_n + \Gamma^2 \beta k \right) \right) \right] \right\} +$$

$$\frac{1}{\beta k} \exp\left(\frac{(\Gamma \beta k)^2}{2} + \beta k(\mu - E_n)\right) \left[\operatorname{erf}\left(\sqrt{\frac{1}{2\Gamma^2}} \left(\mu - E_n + \Gamma^2 \beta k\right)\right) + 1 \right]$$

Evaluating all integration limits, the final form of the free energy \bullet

$$F = \mu N - \frac{eB}{h} \sum_{n,k=1}^{\infty} \left\{ \left[\frac{1}{\beta} \ln(1 + e^{\beta\mu}) \right] \operatorname{erf} \left(\sqrt{\frac{1}{2\Gamma^2}} E_n \right) + \left(\mu - E_n \right) \operatorname{erf} \left(\sqrt{\frac{1}{2\Gamma^2}} (\mu - E_n) \right) + \sqrt{\frac{2\Gamma^2}{\pi}} \left(\exp\left(-\frac{(\mu - E_n)^2}{2\Gamma^2} \right) - \exp\left(-\frac{E_n^2}{2\Gamma^2} \right) \right) + (-1)^k (\mathcal{I}_{nk} - \mathcal{J}_{nk}) \right\}$$
(11)

f(E) is the Fermi-Dirac distribution function.

References

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CONCLUSIONS

- An analytic expression is derived for the free energy of broadened Landau levels of a 2DEG.
- This result can be use to obtain the magnetization and heat capacity closed forms having a Gaussian DOS and compare it with the numerical simulation [4].
- The form can examine the limiting case by deriving the specific heat and comparing the form obtained from Ref. [5].