Tailoring spin orientation via manipulating electronic structure in two-dimensional magnet CrI₃

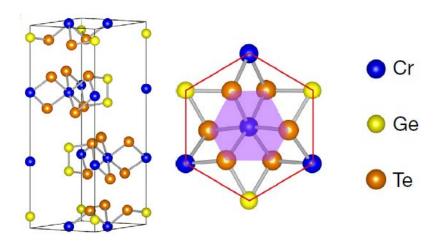
Aug. 28, 2019 at IBS

Jeongwoo Kim

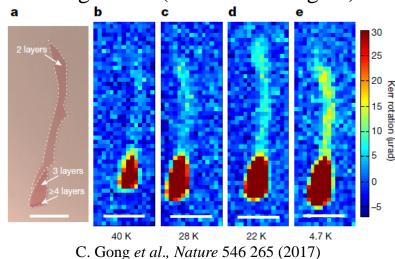
Department of Physics, Incheon National University, Incheon 22012, Korea

Cr₂Ge₂Te₆

Atomic structure

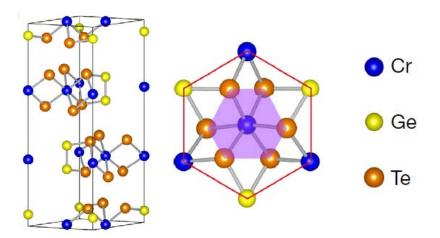


Ferromagnetism (Ker rotation signal)

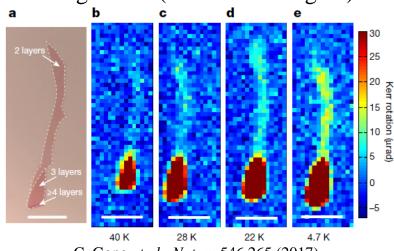




Atomic structure



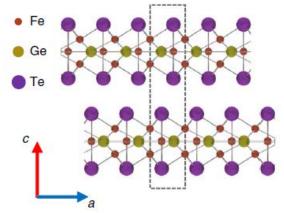
Ferromagnetism (Ker rotation signal)



C. Gong et al., Nature 546 265 (2017)

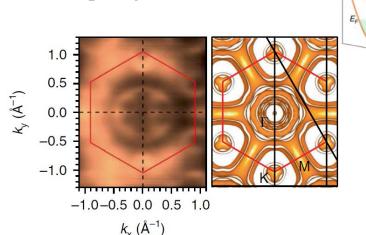
Fe₃GeTe₂

Atomic structure



Z. Fei et al., Nat. Mater. 17 778 (2018)

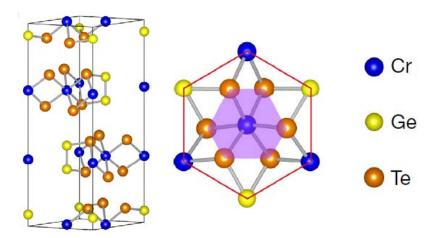
Topological nodal line



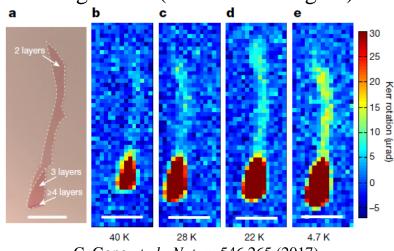
K. Kim et al., Nat. Mater. 17 794 (2018)



Atomic structure



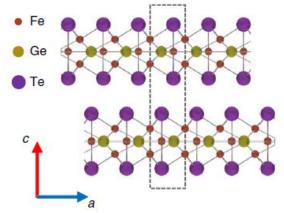
Ferromagnetism (Ker rotation signal)



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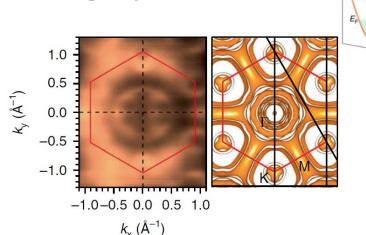
Fe₃GeTe₂

Atomic structure



Z. Fei et al., Nat. Mater. 17 778 (2018)

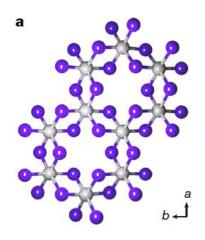
Topological nodal line



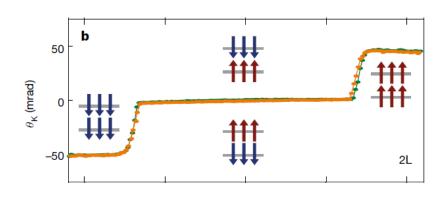
K. Kim et al., Nat. Mater. 17 794 (2018)

CrI₃

Atomic structure

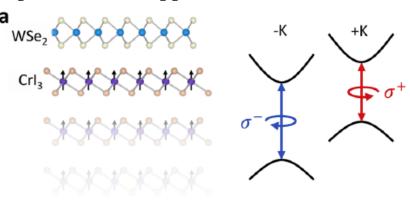


Antiferromagnetic inter-layer coupling



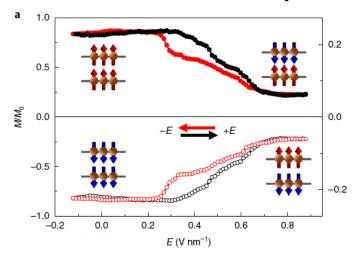
B. Huang et al., Nature 546 270 (2017)

Optoelectronic application (WSe2/CrI3)



K. L. Seyler et al., Nano Lett. 18 3823 (2018)

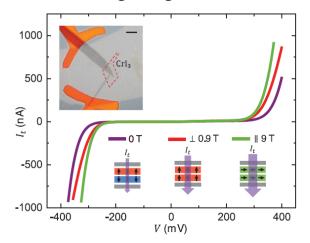
Electrostatic control of the inter-layer coupling

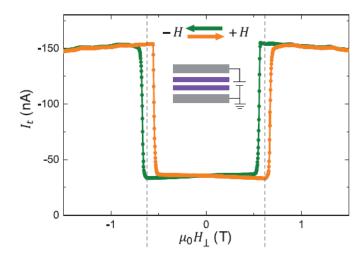


S. Jiang et al., Nat. Mater. 17 406 (2018)

CrI₃

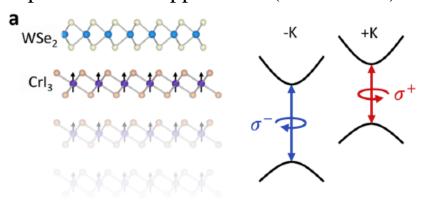
Giant tunneling magnetoresistance





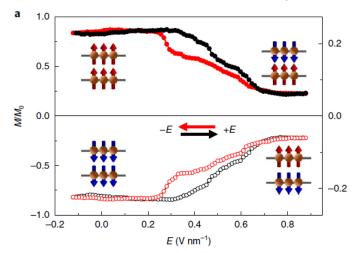
T. Song et al., Science 360 1214 (2018)

Optoelectronic application (WSe2/CrI3)



K. L. Seyler et al., Nano Lett. 18 3823 (2018)

Electrostatic control of the inter-layer coupling



S. Jiang et al., Nat. Mater. 17 406 (2018)

The 2D Heisenberg model, in which the spins can point in any direction, does not have a state with long-range magnetic order.

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The magnetic anisotropy in reduced dimensionality is an essential ingredient for the presence of the 2D magnet.

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The magnetic anisotropy in reduced dimensionality is an essential ingredient for the presence of the 2D magnet.

For physics: an in-depth investigation is prerequisite to the understanding of the underlying physics of the 2D magnetism.

The source of the magnetic anisotropy in CrI₃

- i) Ferromagnetic super-exchange (Major)
- ii) The single ion anisotropy (Minor)

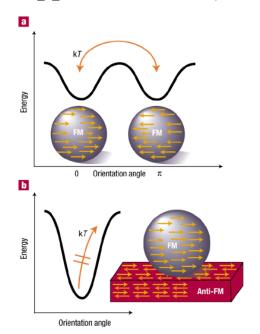
J L Lado et al., 2D Mater. 4 035002 (2017)

The 2D Heisenberg model, in which the spins can point in any direction, does not have a state with long-range magnetic order.



The magnetic anisotropy in reduced dimensionality is an essential ingredient for the presence of the 2D magnet.

For application: the dynamics control and optimization are indispensable



- 1. MAE∝ stability of stored information (Good)
- 2. MAE∝ energy consumption for writing information (Bad)

J Eisenmenger et al., Nat. Mater. 2 437 (2003)

1. What is the electronic origin of the magnetic anisotropy?

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- 2. How can we control the magnetic anisotropy?
 - Dynamically and reversibly
 - The sign and magnitude of MAE

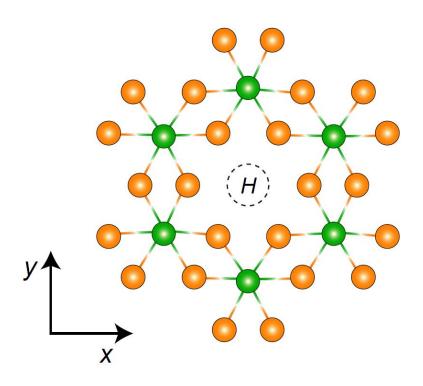
- 1. What is the electronic origin of the magnetic anisotropy?
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 - Dynamically and reversibly
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Enhanced stability and Low-energy operation

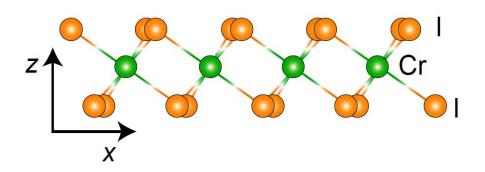
Atomic structure of CrI₃

Top view



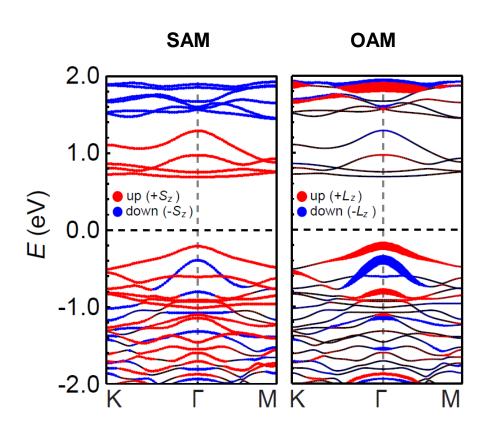
H: hollow site

Side view



D_{3d} point group

Calculated band structure with spin and orbital angular momentum resolution

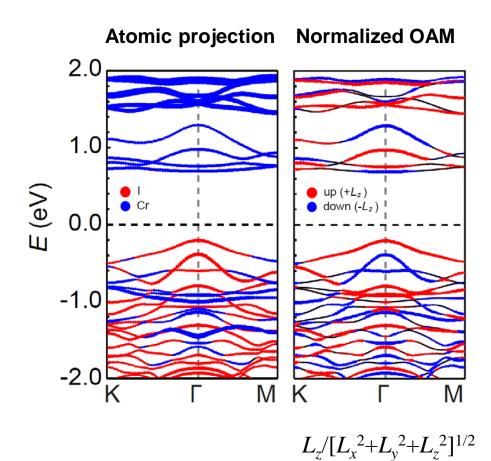


- Cr atom has $3 \mu_{\rm B}$ spin magnetic moment
- Ferromagnetically ordered with the out-of-plane easy axis
- Spin-polarized band structure (left)
- Substantial orbital angular momentum $\langle L_z \rangle$ of two highest valance bands

two highest valance bands (HVBs)

Two lowest conduction bands (LCBs)

Calculated band structure with atomic projection and normalized OAM



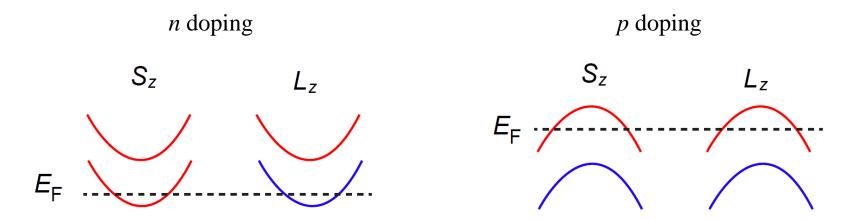
- The orbital angular momentum mainly originates from *p* orbitals in I atoms.
- $\langle L_{\chi} \rangle$, $\langle L_{\gamma} \rangle \ll \langle L_{z} \rangle$
- The discrete rotational symmetry around z axis implies that L_z is a good quantum number.

Evaluation of MAE

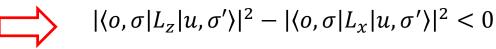
$$MAE = \xi^{2} \sum_{u,\sigma,\sigma'} \sigma \sigma' \frac{|\langle o, \sigma | L_{z} | u, \sigma' \rangle|^{2} - |\langle o, \sigma | L_{x} | u, \sigma' \rangle|^{2}}{E_{u,\sigma} - E_{o,\sigma'}}$$

D.-S. Wang et al., Phys. Rev. B 47 932 (1993)

 ξ : the spin-orbit coupling constant for electrons in Cr u and o: unoccupied and occupied states $E_{u/o,\sigma}$: the energy of the states the spin indices σ and σ '



- The overlap between the two bands with different $\langle L_z \rangle$ values is not significant.
- L_r operator can mix the two HVBs (LCBs)



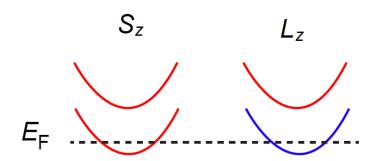
Evaluation of MAE

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D.-S. Wang et al., Phys. Rev. B 47 932 (1993)

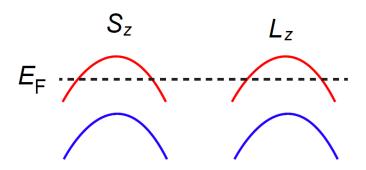
 ξ : the spin-orbit coupling constant for electrons in Cr u and o: unoccupied and occupied states $E_{u/o,\sigma}$: the energy of the states the spin indices σ and σ '

n doping



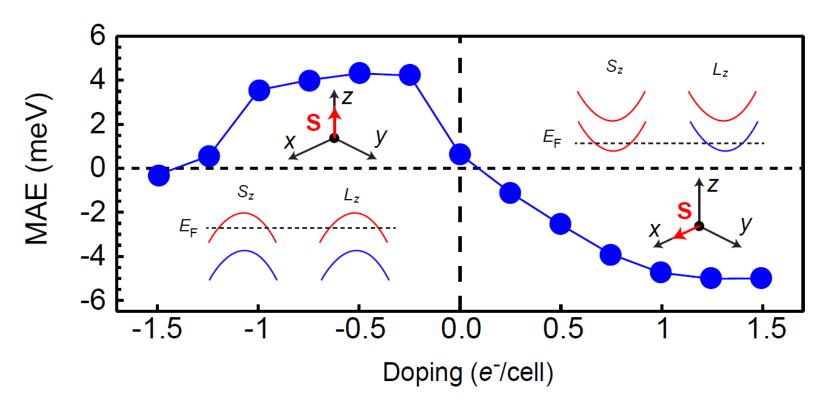
Same spin polarization → negative MAE (in-plane)

p doping



different spin polarization → positive MAE (out-of-plane)

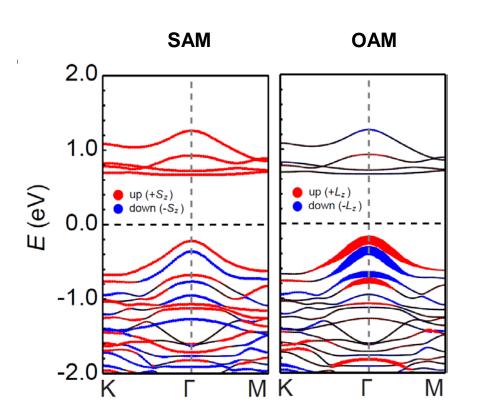
Calculated MAE with various doping levels



- The MAE increases up to 4.35 meV from 0.66 meV ($\sim +560$ % change) by reducing electrons.
- It becomes negative and saturated to -5.01 meV (\sim -860 % change) by adding electrons.

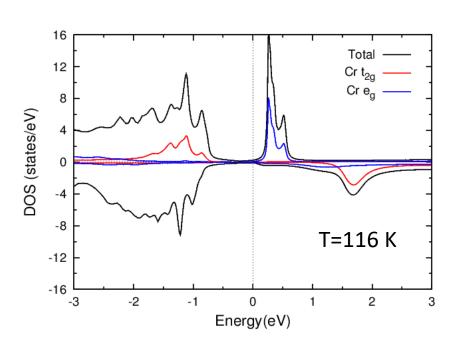
correlation effect

GGA+*U* scheme



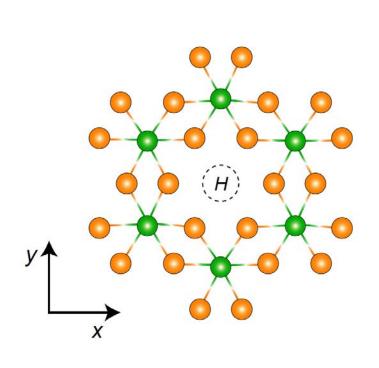
The unique angular momentum distributions of the CrI_3 are retained even though the U effect (2.65 eV) is considered

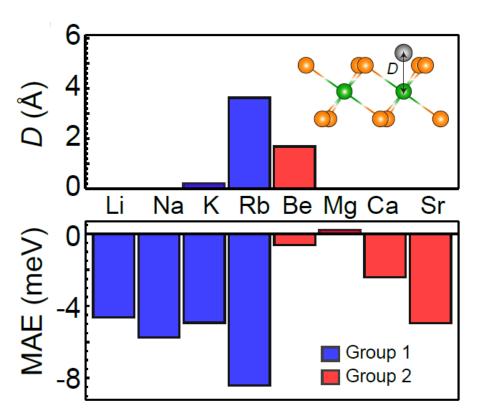
DMFT calculations



- -The hybridization expansion continuous-time quantum Monto Carlo
- Ising-type Coulomb interaction with the on-site Coulomb repulsion U = 10 eV and the Hund's coupling J = 1 eV in Cr 3d orbital

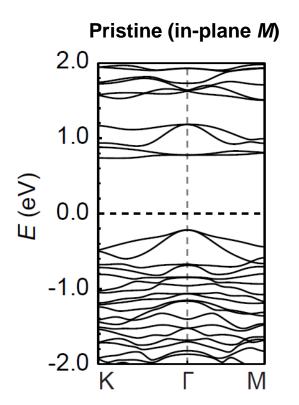
Alkali and Alkaline-earth metal adsorption

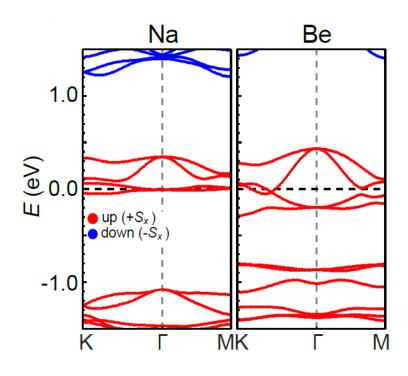




- Alkali metals: yielding sizeable MAEs corresponding to the electron addition case.
- Alkaline-earth metals: deviating from the ideally electron-doped one.

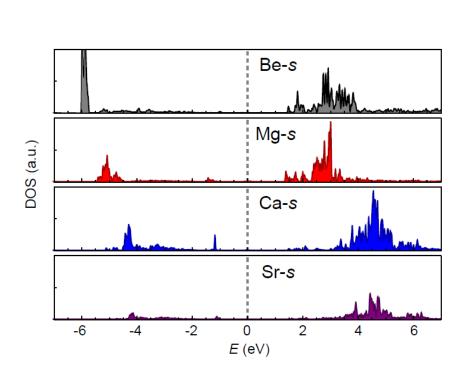
Electronic structures

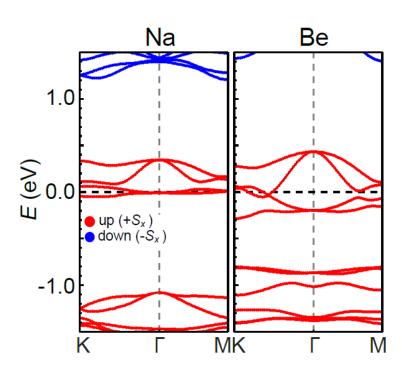




- For the Na adsorption, the band structure remains almost rigid except for the shift of the Fermi level caused by the charge transfer.
- On the other hand, Be atom considerably changes the electronic structure of CrI₃.

Electronic structures



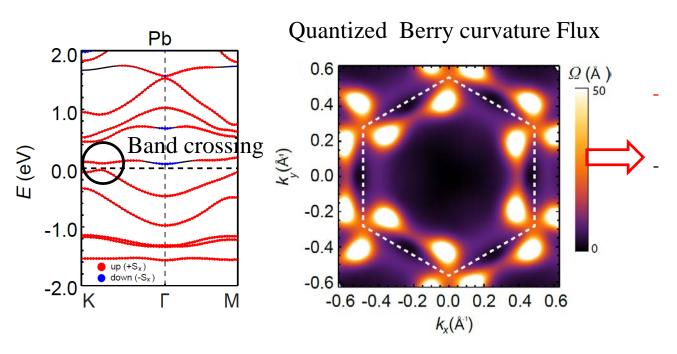


- For the Na adsorption, the band structure remains almost rigid except for the shift of the Fermi level caused by the charge transfer.
- On the other hand, Be atom considerably changes the electronic structure of CrI₃.

Importance of magnetic anisotropy in band topology

The magnetic anisotropy is very sensitive to external perturbations and thus the ground magnetic anisotropy may not be maintained in 2D magnets, sometimes leading to a misleading result without the consideration of the modified spin axis.

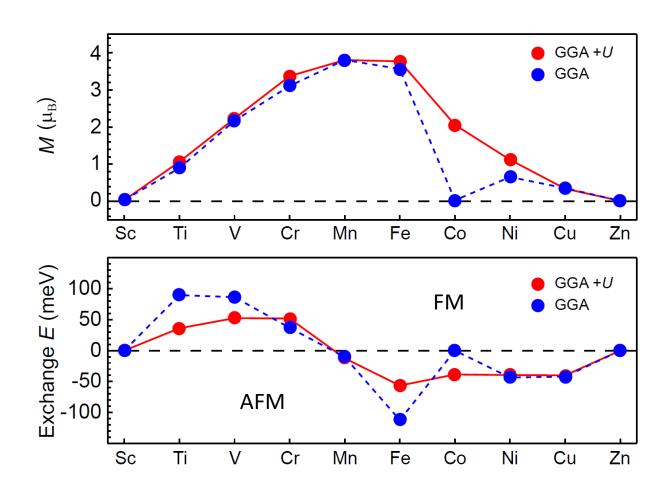
Example: Pb-adsorbed CrI₃



a Chern insulator under the out-of-plane magnetic axis indeed a trivial small-gap insulator with the energetically stable in-plane magnetization

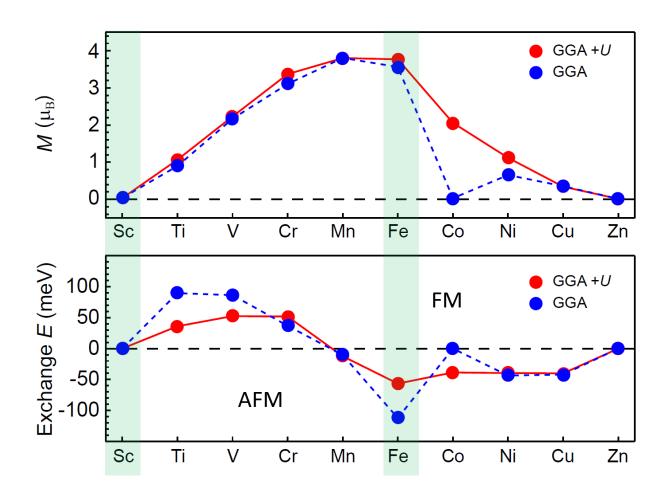
Substitutional doping

Local magnetic moment and exchange energy



Substitutional doping

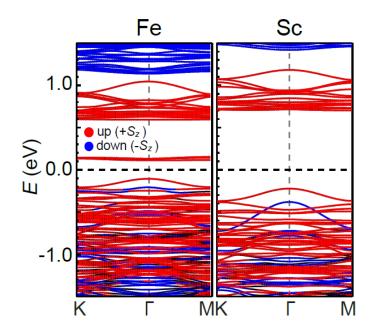
Local magnetic moment and exchange energy



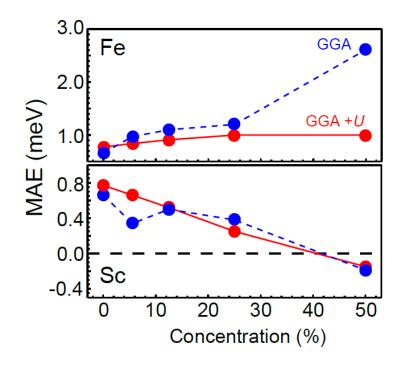
Among 3d transition metals, Fe and Sc atoms maintain the insulating nature of CrI₃ when they are substituted for Cr atoms.

Substitutional doping

Band structure



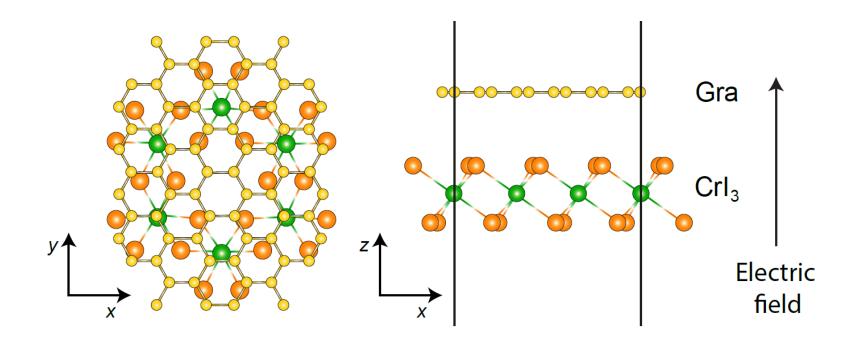
Calculated MAE



- The MAE increases (decreases) with Fe (Sc) concentration.
- the MAE can be effectively adjusted by Fe or Sc doping without compromising the insulating property of CrI₃.

Proximity to graphene

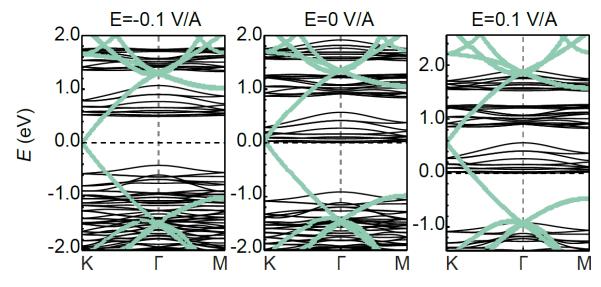
Geometry



a (5×5) graphene is placed on a ($\sqrt{3} \times \sqrt{3}$) CrI₃ monolayer

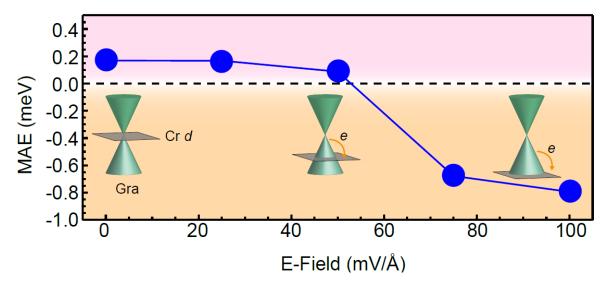
Proximity to graphene

Band structure



- Charge transfer is driven by the perpendicular external field.

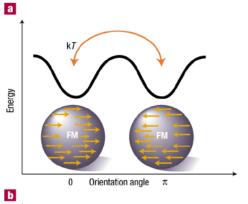
Calculated MAE



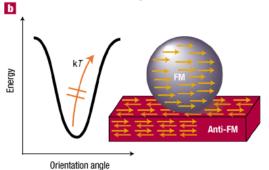
- The spin configuration of CrI_3 can be dynamically controlled by using a perpendicular external field

Proximity to graphene

Dynamical manipulation of MAE



For switching, the energy barrier can be instantaneously vanished by electric field as the magnetic states is reconfigured by nearly zero current.



For stability, the energy barrier is recovered by removing electric field to retain a magnetic state.

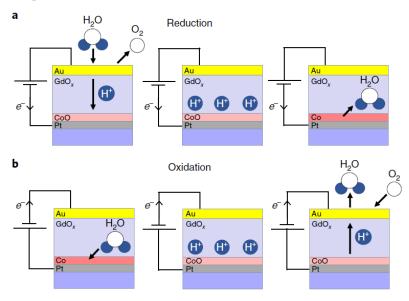
J Eisenmenger et al., Nat. Mater. 2 437 (2003)



energy-efficient spin device

Previous magnetic systems

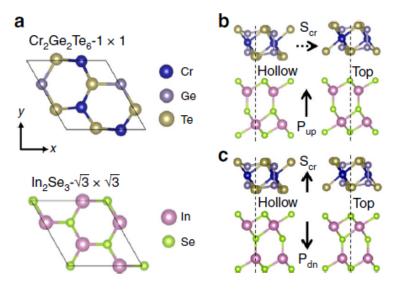
Magneto-ionic control



A. J. Tan et al., Nat. Mater. 18 35 (2019)

- Large variation of MAE, but a long switching time form microseconds to minutes.
- Our charge-transfer mechanism is very fast.

Multiferroicity



C. Gong et al., Nat. Commun. 10 1 (2019)

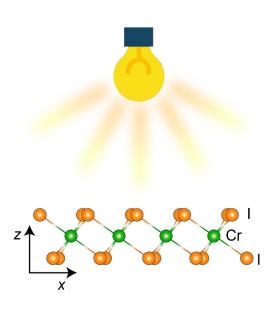
- Difficult to tune the magnitude of MAE
- A wide range of MAE can be accessible by adjusting electric field in our case.



High-speed and reversible spin device

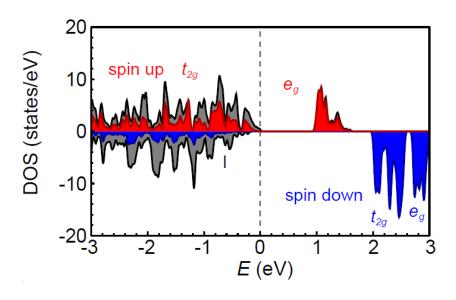
Optical excitation

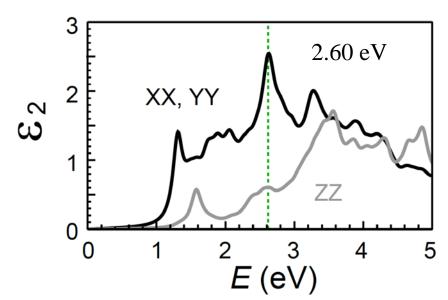
Carrier excitation through irradiation



We consider the manipulation of the MAE of a pristine and stand-alone CrI₃ monolayer through electronic excitation obtained by using a resonant laser pulse.

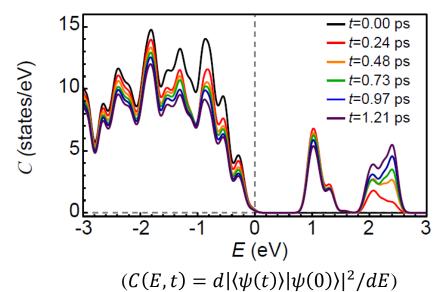
Choice of Light frequency

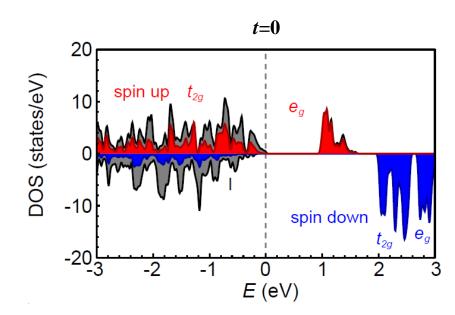




Optical excitation

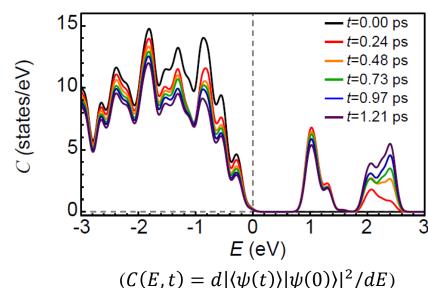
Time-varying DOS of occupied states





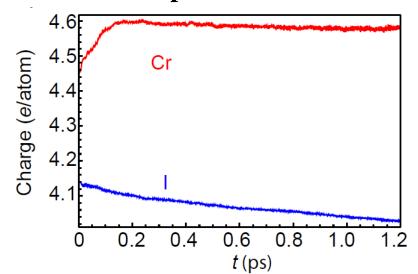
Optical excitation

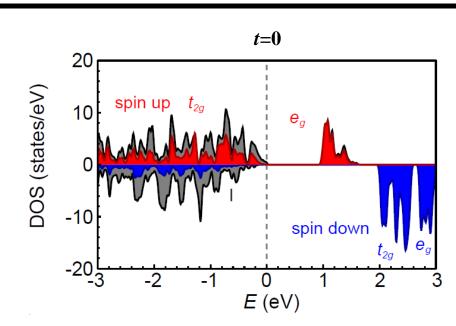
Time-varying DOS of occupied states

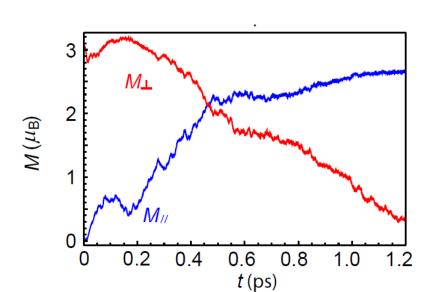


$(C(E,t) = d|\langle \psi(t)\rangle|\psi(0)\rangle|^2/dE)$

Reorientation of spins







Summary

1. We performed first-principles calculations for MAE in ${\rm CrI_3}$ -based 2D structures and presented a few accessible methods to manipulate its magnetization configurations.

2. As a result of extensive DFT calculation, we found that, in the determination of the magnetization configuration, the spin and orbital angular momentums between the occupied and unoccupied bands are intertwined, giving rise to the drastic variation of the MAE upon the band filling

3. Especially, we demonstrated that the anisotropic magnetic configuration can be easily altered by the perpendicular external field, implying high potential for energy-efficient fast switching of two-dimensional magnets with nearly zero critical current density.