Topological magnetoelastic excitations in triangular antiferromagnets

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Motivation: can non-zero Chern number arise from magnon-phonon coupling?

- Topological magnon (R. Shindou et al PRB 2013) and topological phonon (L. Zhang et al PRL 2010) bands with non-zero Chern number implies existence of chiral edge states and non-zero thermal Hall conductivity.
- Magnon and phonon can naturally hybridize due to magnon-phonon coupling, with large Berry curvature at anticrossing region (Takahashi, Nagaosa PRL 2016)
- We give a toy model on triangular antiferromagnet where topologically trivial magnon and



phonon hybridize to form topologically nontrivial magnetoelastic bands. Magnon-phonon interaction can modify the thermal Hall conductivity.



Phonon

- $\mathcal{H}_p = \frac{1}{2} \sum \left[\{ \boldsymbol{p}_{\alpha}(\boldsymbol{R})^2 + 2\boldsymbol{u}_{\alpha}(\boldsymbol{R})A_{\alpha\alpha}\boldsymbol{p}_{\alpha}(\boldsymbol{R}) \} \delta_{\alpha\beta}\delta_{\boldsymbol{R},\boldsymbol{R}'} + \boldsymbol{u}_{\alpha}(\boldsymbol{R}) \{ K_{\alpha\beta}(\boldsymbol{R}-\boldsymbol{R}') (A^2)_{\alpha\beta} \} \boldsymbol{u}_{\beta}(\boldsymbol{R}') \right].$
 - Nearest Neighbor spring constant and Phenomenological Raman coupling

h = 0

(a)





- Although the magnon-phonon Hamiltonian is often written using the fields



(b)

Magnon-phonon coupling

- $\mathcal{H}_c = \sum K_{mp} \boldsymbol{R}_{ij} \cdot \Delta \boldsymbol{u}_{ij} \boldsymbol{S}_i \cdot \boldsymbol{S}_j$
- Exchange striction: $J(|\mathbf{r}_i \mathbf{r}_j|) \approx J + K_{mp} \mathbf{R}_{ij} \cdot \Delta \mathbf{u}_{ij}$,
- where $r_i = R_i + u_i$





Fig. Caption (a), (c), (e): Berry curvatures calculated by using field Ψ_k (b), (d), (f): Berry curvatures calculated by using field Σ_k

- $\Psi_{k} = (a_{k}, b_{1,k}, b_{2,k}, a_{-k}^{\dagger}, b_{1,-k}^{\dagger}, b_{2,-k}^{\dagger}),$
- where a's are Holstein Primakoff operators and b are phonon operators, this does not give the correct Berry curvature.
- We should write the Hamiltonian using

 $\Phi_{k} = \left(a_{k}, a_{-k}^{\dagger}, \boldsymbol{p}_{k}^{T}, \boldsymbol{u}_{k}^{T}\right)$

Then, $\Sigma_{k} \equiv PV^{\dagger}\Phi_{k}$ is bosonic BdG field: •



Thermal Hall conductivity



Fig. Caption:

Left: thermal Hall conductivity computed using Berry curvatures calculated by using field Σ_k Right: thermal Hall conductivity computed using Berry curvatures calculated by using field Ψ_k

Note that the phonon Berry curvature is missed in (a), (c), and (e).

The Hamiltonian written using Ψ_k does not have to be smooth, which results in unphysical singularities in Berry curvature.

The Chern numbers for (b), (d), and (f) are -2, 4, -2 respectively.

Spin Current?

- Consider a collinear antiferromagnet.
- The magnon bands carry spin quantum numbers
- The magnon-phonon interaction can induce thermal Hall and spin Nernst current





Fig. Caption

Top: Thermal Hall and spin Nernst conductivity Bottom: spin density

Nontrivial bulk boundary correspondence between the spin Nernst current and the edge spin accumulation

Thermal gradient induces (bulk) spin density, which is allowed because magnon-phonon interaction violates spin conservation.