

# Polariton rings and ring- based metamaterials

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# Outline

- a. Magnetic field in classical dynamics
- b. Magnetic field in quantum physics and AB effect
- c. Optically induced AB effect in mesoscopic rings
- d. Systems of interconnected rings
- e. Optical AB effect for excitons
- f. Polariton rings and geometric phase
- g. Edge states in polariton ring lattices

# Regimes of light- matter interaction

2DEG irradiated by electromagnetic field



Weak light-matter interaction

Strong light-matter interaction

Electron spectrum unperturbed

Electron spectrum modified  
(dynamic Stark effect)

Interact via photon  
emission/absorption

Interact via electron “dressing”

Effects: photovoltaic, high-  
frequency conductivity

Effects: **exciton-polaritons in microcavities**; band gap opening and metal-insulator transition in graphene; **artificial U(1) gauge fields**

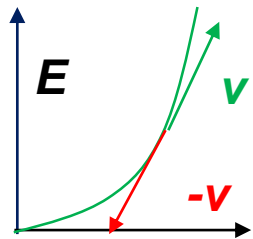
# Magnetic field in classical and quantum physics

## Classical physics

### Lorentz Force

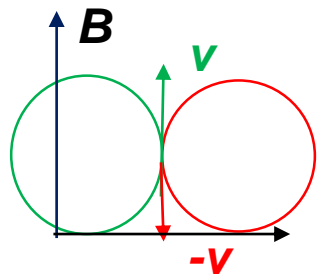
$$\frac{d^2 \mathbf{r}}{dt^2} = \mathbf{F} = q\mathbf{E} + q[\mathbf{v} \times \mathbf{B}]$$

### Time inversion symmetry



$$t \rightarrow -t$$

$$E \rightarrow E$$



$$t \rightarrow -t$$

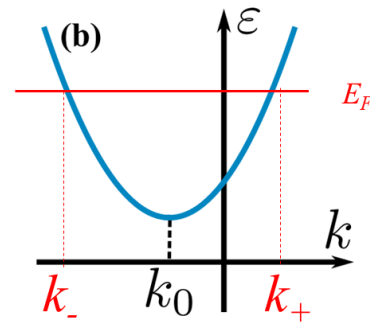
$$B \rightarrow -B$$

## Quantum physics

$$\hat{H}_0 = \frac{1}{2m_e} (\hat{p}_\varphi - eA_\varphi)^2 \quad \hat{p}_\varphi = i\hbar \frac{\partial}{\partial \varphi} \quad A_\varphi = \frac{\Phi}{2\pi R}$$

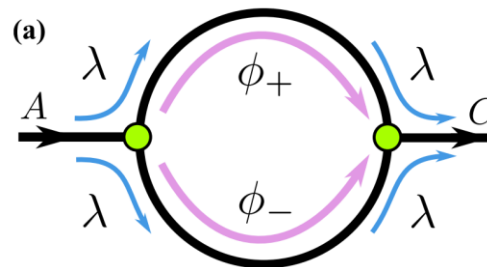
$$\Psi_m = \frac{1}{\sqrt{2\pi}} e^{im\varphi} \quad k = m/R$$

$$\varepsilon(m) = \frac{\hbar^2}{2m_e R^2} \left(m + \frac{\Phi}{\Phi_0}\right)^2 \quad \varepsilon(k) = \frac{\hbar^2}{2m_e} \left(k + \frac{\Phi}{R\Phi_0}\right)^2$$

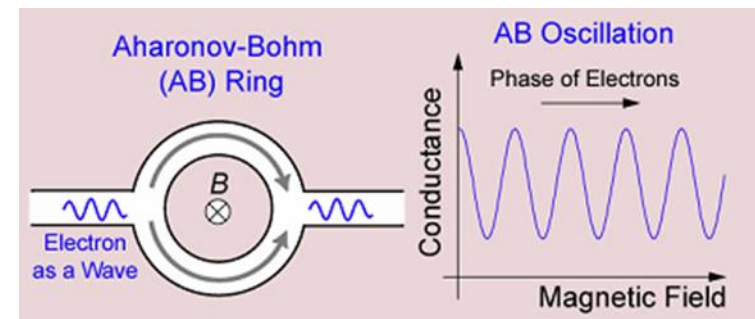


### Time reversal asymmetry

$$E(k) \neq E(-k) \longrightarrow \Delta\phi = \pi R(k_+ - k_-) \neq 0$$



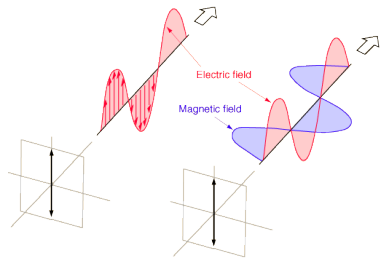
$$G = 2 \frac{e^2}{h} \left( 1 - \left| \frac{\sin^2\left(\frac{\Phi}{2\Phi_0}\right)}{1 - \exp(2\pi i k_F R) \cos^2\left(\frac{\Phi}{2\Phi_0}\right)} \right|^2 \right)$$



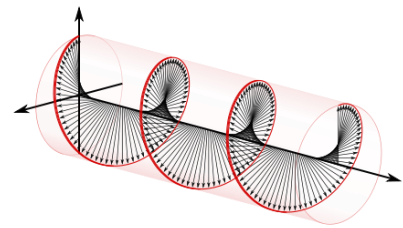
Aharonov-Bohm oscillations in conductance

# Time- reversal symmetry breaking by circular polarized light

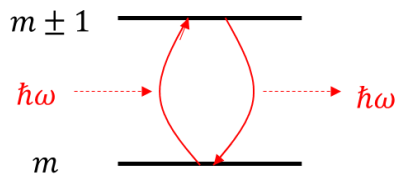
AB effect is a consequence of non-equivalency of clockwise and anticlockwise motion, i.e. of the breaking of time- reversion symmetry



Linear polarized light: time inversion does not change polarization



Circular polarized light: time inversion changes polarization



$$\hat{H}_0 = \frac{1}{2m_e} (\hat{p}_\varphi - eA_\varphi)^2 \quad \hat{\mathcal{H}} = \hat{\mathcal{H}}_R^{(0)} + \hat{U}_R$$

$$\hat{U}_R = -ieR\sqrt{\frac{\pi\hbar\omega}{V}}(e^{i\varphi}\hat{a} - e^{-i\varphi}\hat{a}^\dagger) \quad \varepsilon_{m,N} = \varepsilon_{m,N}^{(0)} + \frac{|\langle m+1, N-1 | \hat{U}_R | m, N \rangle|^2}{\varepsilon_{m,N}^{(0)} - \varepsilon_{m+1, N-1}^{(0)}} + \frac{|\langle m-1, N+1 | \hat{U}_R | m, N \rangle|^2}{\varepsilon_{m,N}^{(0)} - \varepsilon_{m-1, N+1}^{(0)}}$$

$$\varepsilon(m) = m^2 \varepsilon_R + eE_0 R \left[ \frac{eE_0 R / 2\varepsilon_R}{(2m - \hbar\omega / \varepsilon_R)^2 - 1} \right]$$

$$\varepsilon(k) = \frac{\hbar^2 k^2}{2me} + \frac{\hbar e^2 E_0^2}{2me^2 R \omega^3} k = \frac{\hbar^2}{2me} \left( k + \frac{E_0^2}{2me\hbar 2R\omega^3} \right)^2 + const$$

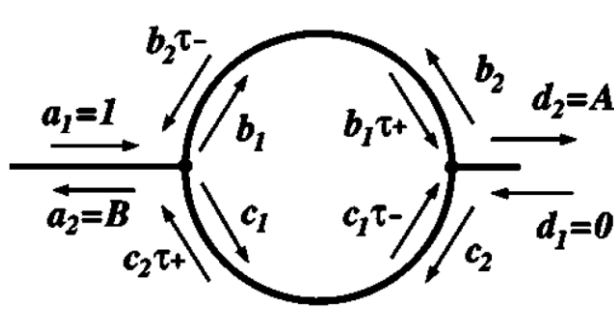
$$\varepsilon(k) = \frac{\hbar^2}{2me} \left( k + \frac{\Phi}{R\Phi_0} \right)^2$$

Illumination of the ring by circular polarized light leads to the appearance of an effective magnetic flux and syntetic U(1) gauge potential

$$\Phi_{eff} = \frac{e\pi E_0^2}{2me\omega^3}$$

$$A_{eff} = \frac{eE_0^2}{4me\omega^3 R}$$

# Optical AB oscillations



$$\begin{pmatrix} b_1 \\ a_2 \\ c_1 \end{pmatrix} = \begin{pmatrix} r & \varepsilon & t \\ \varepsilon & \sigma & \varepsilon \\ t & \varepsilon & r \end{pmatrix} \cdot \begin{pmatrix} b_2 \tau_- \\ 1 \\ c_2 \tau_+ \end{pmatrix}$$

$$\begin{pmatrix} b_2 \\ d_2 \\ c_2 \end{pmatrix} = \begin{pmatrix} r & \varepsilon & t \\ \varepsilon & \sigma & \varepsilon \\ t & \varepsilon & r \end{pmatrix} \cdot \begin{pmatrix} b_1 \tau_+ \\ 0 \\ c_1 \tau_- \end{pmatrix}$$

$$\tau_{\pm} = e^{i\pi k_{\pm} R}$$

$$k_{\pm} = k \pm \frac{\Phi_{eff}}{2\Phi_0}$$

$$G = \frac{2e^2}{h} |A|^2 \quad \text{For fully transparent contacts } (\sigma=0):$$

$$G = \frac{2e^2}{h} \left[ 1 - \left| \frac{\sin^2(\Phi_{eff}/2\Phi_0)}{1 - \exp(i2\pi R k_F) \cos^2(\Phi_{eff}/2\Phi_0)} \right|^2 \right]$$

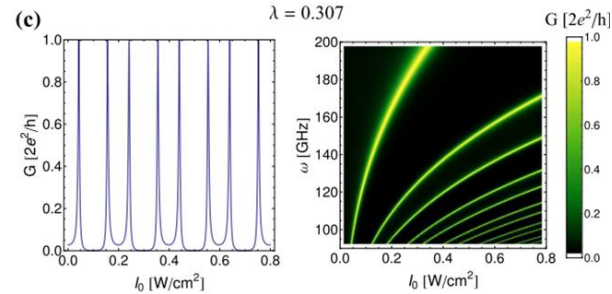
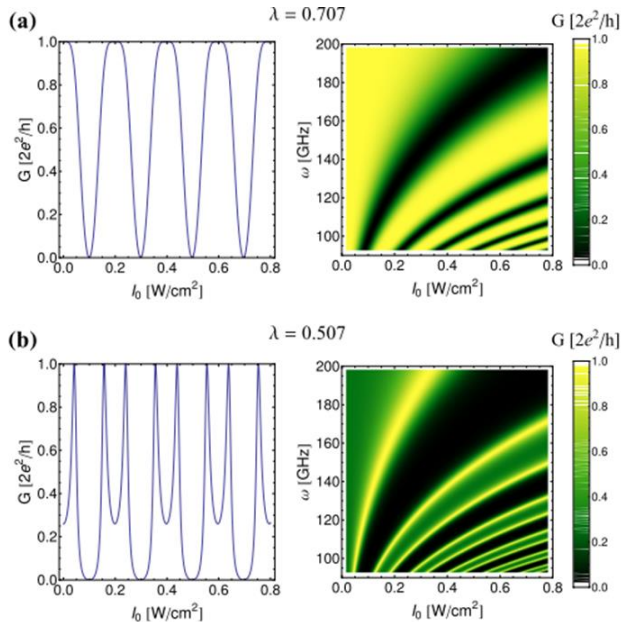
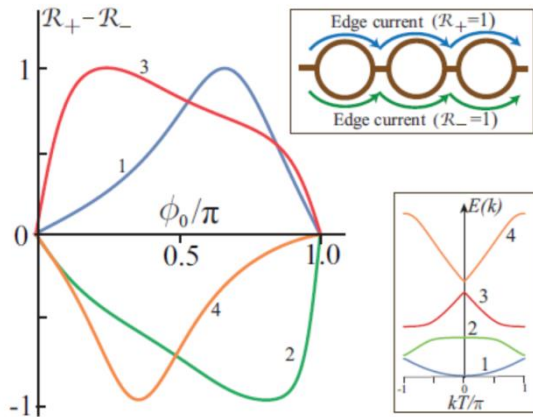
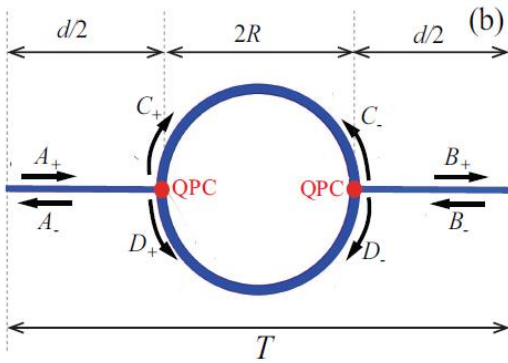
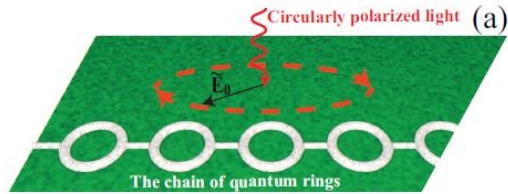


FIG. 2: (Color online) Conductance of a mesoscopic ring,  $G$ , under a circularly polarized electromagnetic wave as a function of wave intensity  $I_0$  and wave frequency  $\omega$ . Plots (a), (b) and (c) correspond to different transmission amplitudes  $\lambda$  between the current leads and the ring. Frames in the left column are fixed at the wave frequency  $\omega = 100$  GHz. In all plots, the ring parameters are assumed to be  $R = 10 \mu m$ ,  $\varepsilon_F = 10$  meV, and  $m_e = 0.1 m_{e0}$ , where  $m_{e0}$  is the mass of free electron.

# 1D array of rings

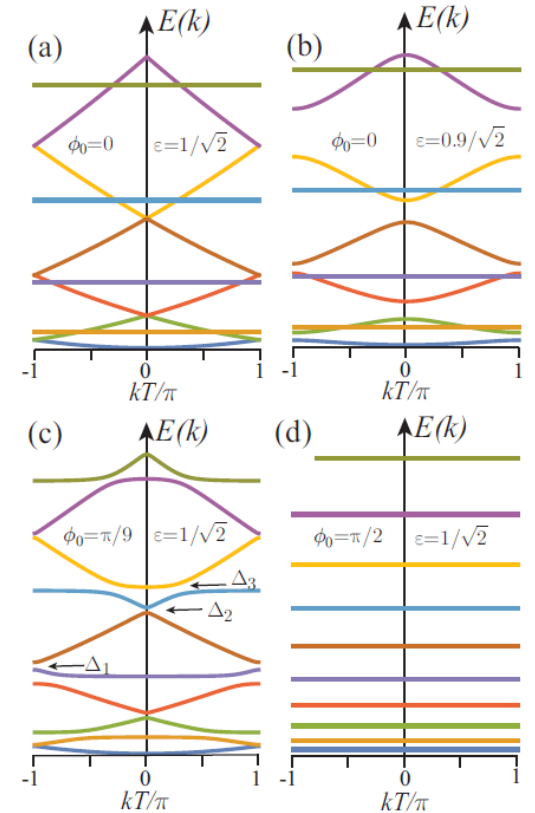


$$\begin{pmatrix} A_+ \\ A_- \end{pmatrix} = e^{ikT} \begin{pmatrix} B_+ \\ B_- \end{pmatrix}$$

$$\begin{pmatrix} A_- e^{-iqd/2} \\ C_+ \\ D_+ \end{pmatrix} = S \begin{pmatrix} A_+ e^{iqd/2} \\ C_- e^{i(\pi qR - \phi_0)} \\ D_- e^{i(\pi qR + \phi_0)} \end{pmatrix},$$

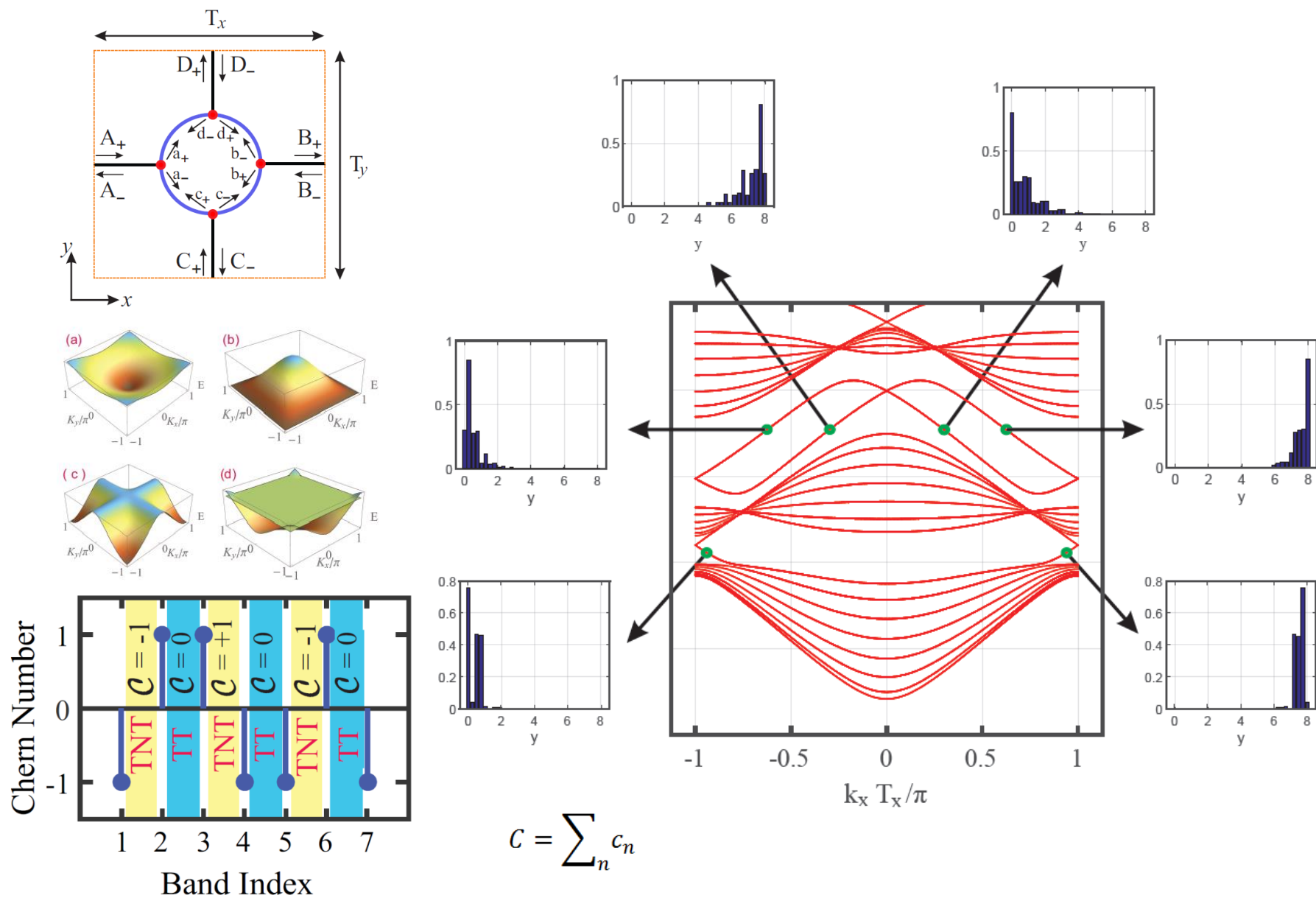
$$\begin{pmatrix} B_+ e^{-iqd/2} \\ C_- \\ D_- \end{pmatrix} = S \begin{pmatrix} B_- e^{iqd/2} \\ C_+ e^{i(\pi qR + \phi_0)} \\ D_+ e^{i(\pi qR - \phi_0)} \end{pmatrix},$$

$$S = \begin{pmatrix} \sqrt{1-2\varepsilon^2} & \varepsilon & \varepsilon \\ \varepsilon & \frac{-(1+\sqrt{1-2\varepsilon^2})}{2} & \frac{(1-\sqrt{1-2\varepsilon^2})}{2} \\ \varepsilon & \frac{(1-\sqrt{1-2\varepsilon^2})}{2} & \frac{-(1+\sqrt{1-2\varepsilon^2})}{2} \end{pmatrix}$$



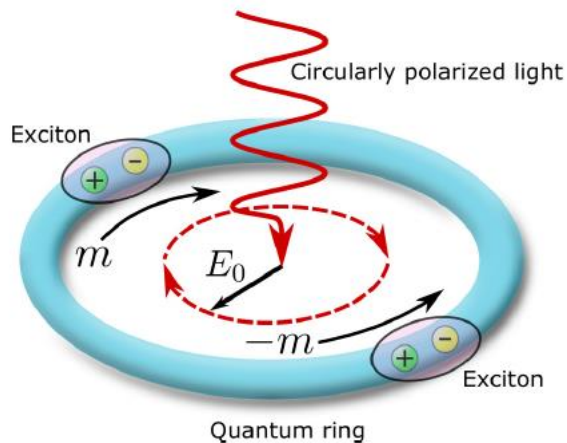
$$\begin{aligned} & \sin(qd) \left[ (1 - \varepsilon^2) \cos^2 \phi_0 + \sqrt{1 - 2\varepsilon^2} \sin^2 \phi_0 \right. \\ & \left. - \cos(2\pi qR) \right] + \varepsilon^2 \left[ \sin[q(d - 2\pi R)] \right. \\ & \left. + 2 \cos \phi_0 \sin(\pi qR) \cos(kT) \right] = 0, \end{aligned}$$

# 2D array of rings





# Optical AB effect for excitons



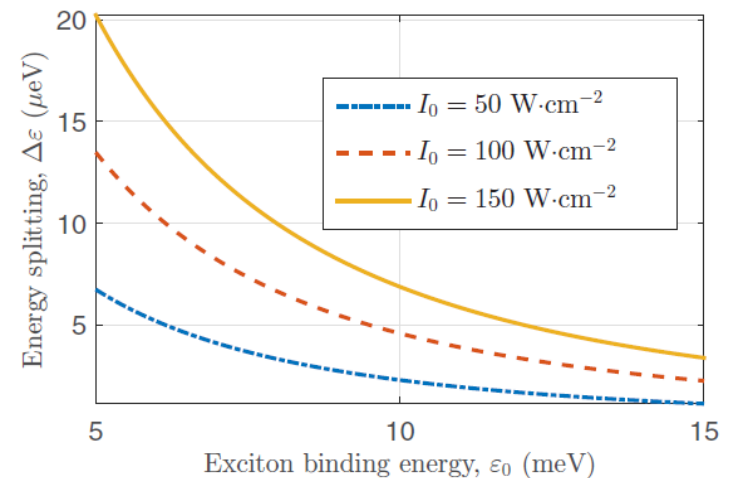
$$\hat{U} = \frac{iq_e R}{2} \sqrt{\frac{\hbar\omega}{\epsilon_0 V_0}} [(e^{-i\varphi_e} - e^{-i\varphi_h})\hat{a}^\dagger + (e^{i\varphi_h} - e^{i\varphi_e})\hat{a}]$$

$$\Delta\tilde{\epsilon}_{0,m} = \frac{\hbar\omega}{2} \left( \frac{m_h - m_e}{M} \right)^2 (eE_0 a)^2 \left[ \frac{1}{\epsilon_R^2 (1 - 2m)^2 - (\hbar\omega)^2} - \frac{1}{\epsilon_R^2 (1 + 2m)^2 - (\hbar\omega)^2} \right]. \quad (12)$$

Exciton Bohr radius

$$\Phi_{eff} = \frac{e\pi E_0^2}{2me\omega^3}$$

Effective flux depends on mass! And masses of an electron and a hole are different, so there appears uncompensated flux for the exciton!



# Geometric Berry Phase

$$H(x(t))|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

$$H(x(t))|n(x(t))\rangle = E_n(x(t))|n(x(t))\rangle$$

**Adiabatic approximation:**

$$|\psi(t)\rangle = e^{i\phi_n} |n(x(t))\rangle$$

$$\phi_n = \theta_n + \gamma_n$$

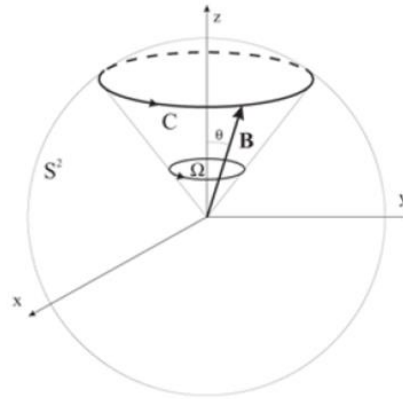
**Dynamic phase**

$$\theta_n(t) = -\frac{1}{\hbar} \int_0^t E_n(\tau) d\tau$$

**Geometric phase**

$$\gamma_n(C) = i \oint_C \langle n(x) | \nabla | n(x) \rangle dx$$

**Geometric phase:** adiabatic evolution of spin in slowly changing time dependent magnetic field



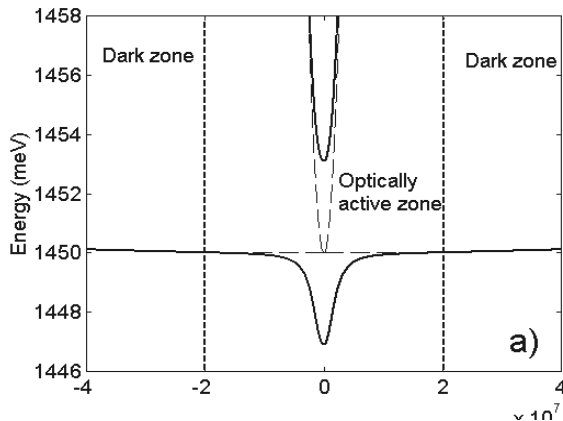
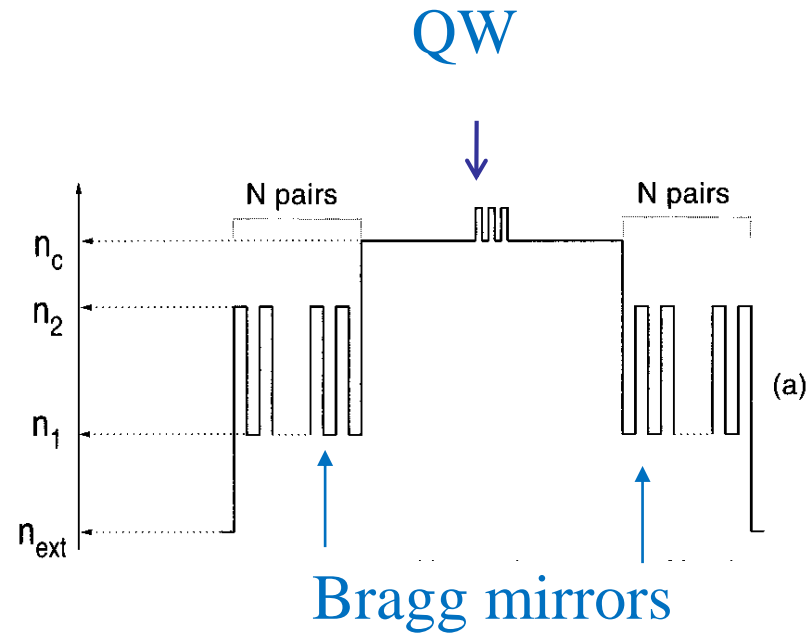
$$\vec{B}(t) = B_0 \begin{pmatrix} \sin \theta \cos(\omega t) \\ \sin \theta \sin(\omega t) \\ \cos \theta \end{pmatrix}$$

$$|n_+(t)\rangle = \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\omega t} \sin \frac{\theta}{2} \end{pmatrix}, \quad |n_-(t)\rangle = \begin{pmatrix} -\sin \frac{\theta}{2} \\ e^{i\omega t} \cos \frac{\theta}{2} \end{pmatrix}$$

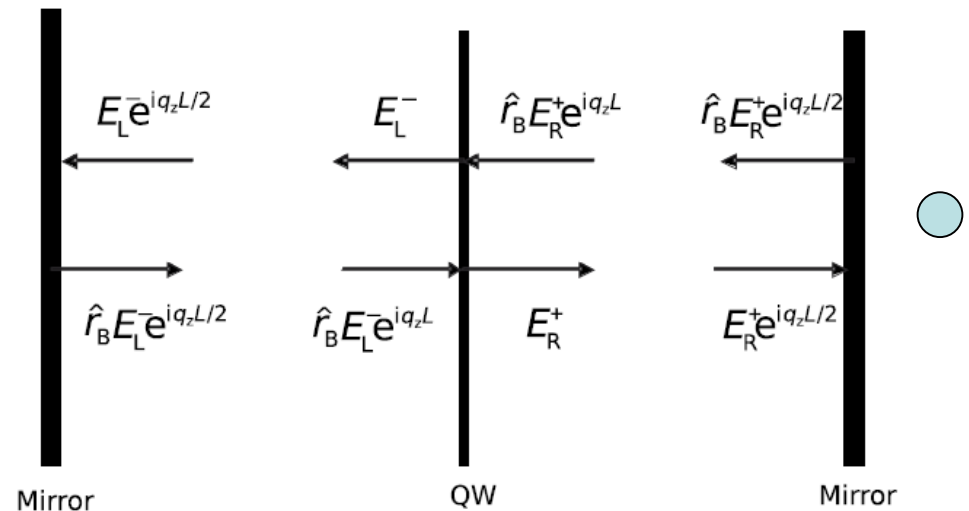
$$\gamma_{\pm}(C) = -\pi(1 \mp \cos \theta)$$

$$\gamma_{\pm}(C) = \mp \frac{1}{2} \Omega(C)$$

# Quantum microcavities



The role of the Bragg mirrors: confinement of electromagnetic field. For infinite number of the layers the photonic bandgap is formed in the mirrors, and photon becomes perfectly confined inside a cavity.

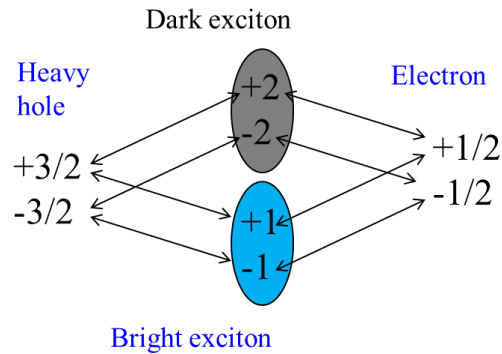


Confined photon acquires a mass

$$\hbar \omega_c = \hbar \frac{c}{n_c} \sqrt{q^2 + k_z^2} \approx \hbar \frac{c}{n_c} k_z \left( 1 + \frac{q^2}{2k_z^2} \right) \equiv \frac{\hbar c}{n_c L_c} + \frac{\hbar^2 q^2}{2m_{ph}}$$

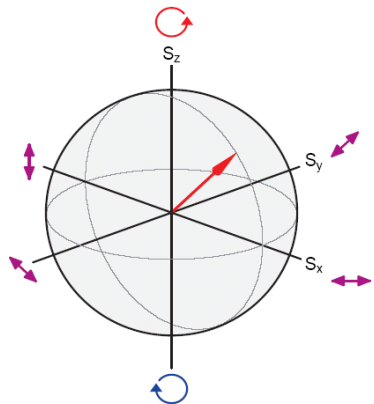
$$m_{ph} = \frac{\hbar n_c}{c L_c} \sim 10^{-4} \div 10^{-5} m_e$$

# Spin and polarization of cavity polaritons



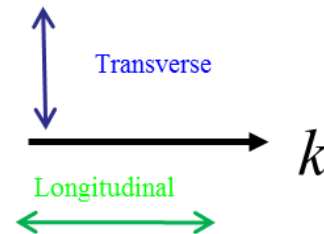
## Density matrix of a two level system

$$\rho_k^- = \frac{N_k^-}{2} I + \vec{S}_k^- \cdot \vec{\sigma} = \begin{pmatrix} N_k^- / 2 + S_{k,z}^- & S_{k,x}^- - i S_{k,y}^- \\ S_{k,x}^- + i S_{k,y}^- & N_k^- / 2 - S_{k,z}^- \end{pmatrix}$$



$$S_z = \frac{1}{2} (N_{\uparrow} - N_{\downarrow})$$

## Polariton TE- TM splitting

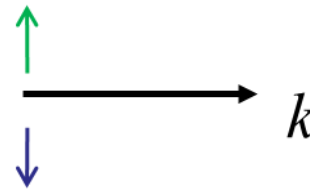


$$H = \begin{pmatrix} H_0(k) & \beta(k_x - ik_y) \\ \beta(k_x + ik_y) & H_0(k) \end{pmatrix} = H_0(k) \mathbf{I} + \Omega_{LT} \cdot \vec{\sigma}$$

$$\Omega_{LT} = |\Omega_{LT}| (\mathbf{e}_x \sin 2\varphi + \mathbf{e}_y \cos 2\varphi)$$

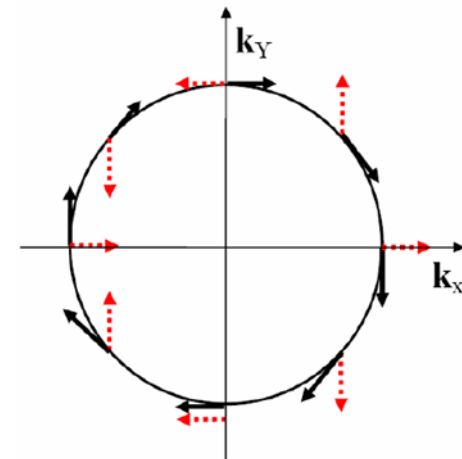
$$\vec{\Omega}_{LT} = \vec{\Omega}_{LT,ex} |X_E^2| + \vec{\Omega}_{LT,ph} |X_P^2|$$

## Analog for electrons: Rashba splitting

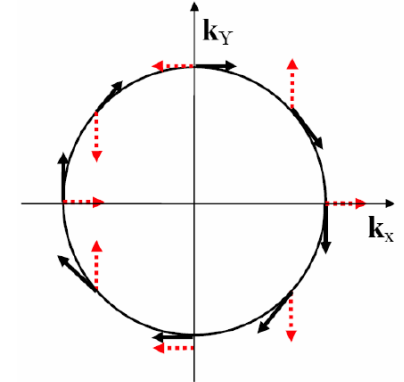
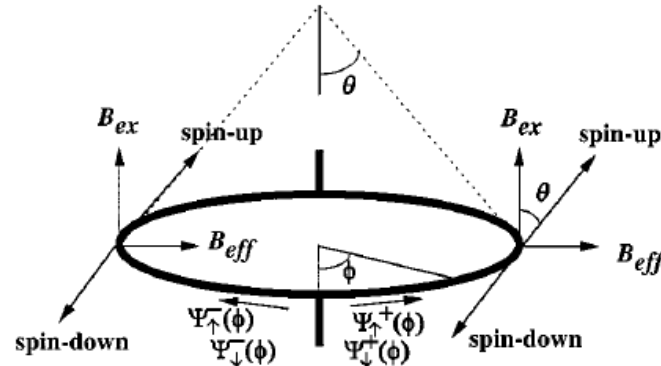
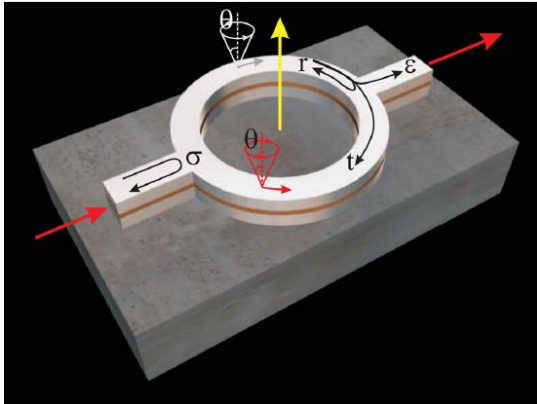


$$H = \begin{pmatrix} H_0(k) & \alpha(k_x - ik_y) \\ \alpha(k_x + ik_y) & H_0(k) \end{pmatrix} = H_0(k) \mathbf{I} + \Omega_R \cdot \vec{\sigma}$$

$$\Omega_R = |\Omega_R| (\mathbf{e}_x \sin \varphi + \mathbf{e}_y \cos \varphi)$$



# Optical analog of AB interferometer



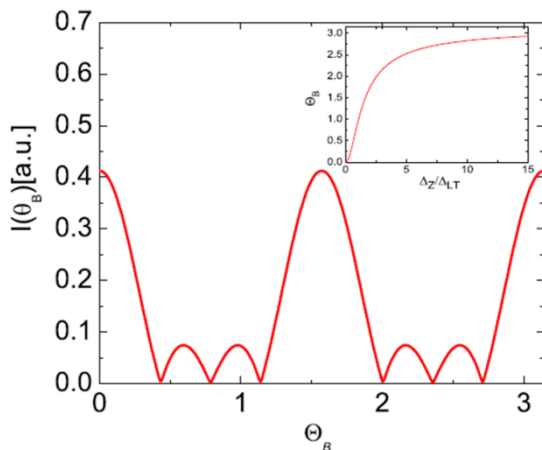
$$H = \begin{pmatrix} -\frac{\hbar^2}{2mR^2} \frac{\partial^2}{\partial \varphi^2} - \frac{\Delta_Z}{2} & \frac{1}{2} \Delta_{LT} e^{2i\varphi} \\ \frac{1}{2} \Delta_{LT} e^{-2i\varphi} & -\frac{\hbar^2}{2mR^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\Delta_Z}{2} \end{pmatrix}$$

← TE-TM splitting  
← Zeeman splitting

$$\Psi_+(\phi) = \frac{1}{\sqrt{1+\xi^2}} \begin{pmatrix} -\xi e^{+i\phi} \\ e^{-i\phi} \end{pmatrix} e^{ik_+ a \phi}$$

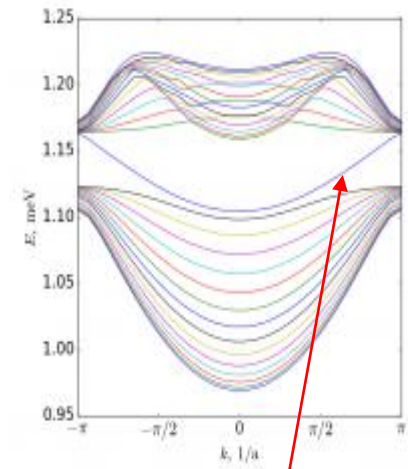
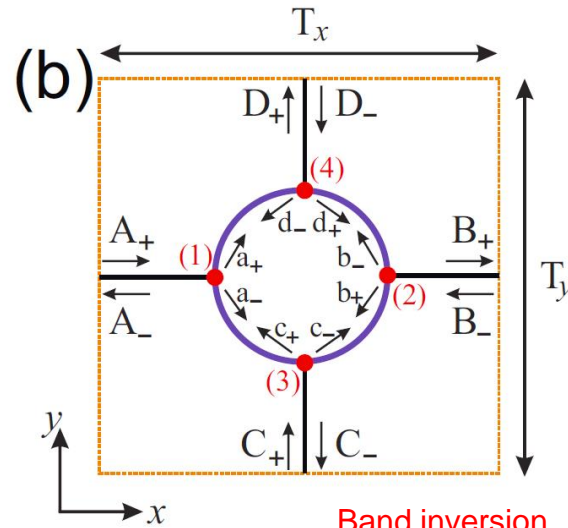
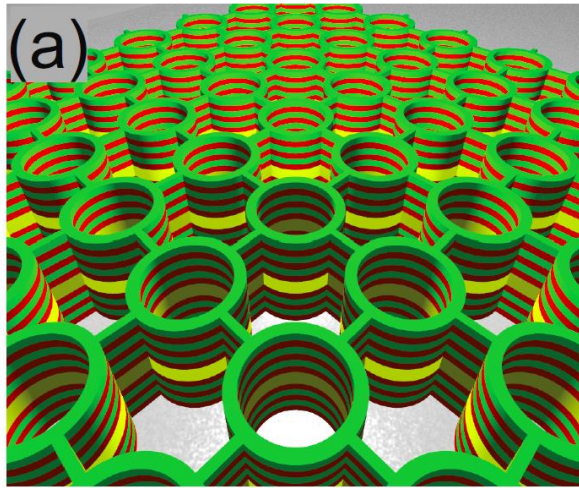
$$\Psi_-(\phi) = \frac{1}{\sqrt{1+\xi^2}} \begin{pmatrix} e^{+i\phi} \\ \xi e^{-i\phi} \end{pmatrix} e^{ik_- a \phi}$$

$$\xi = \frac{\Delta_{LT}/2\Delta_Z}{1 + \sqrt{(\Delta_{LT}/2\Delta_Z)^2 + 1}}$$

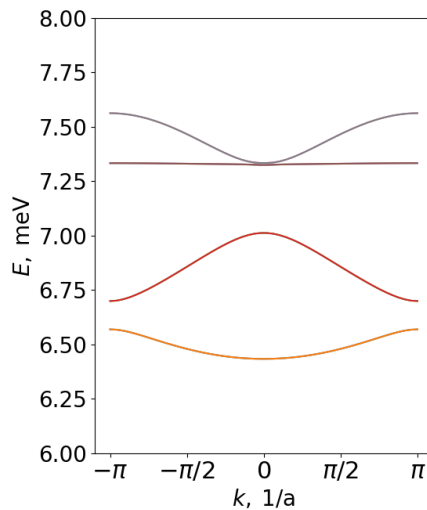


$$\theta_B = \pm \pi \left( 1 - \frac{\Delta_Z}{\sqrt{\Delta_Z^2 + \Delta_{LT}^2}} \right)$$

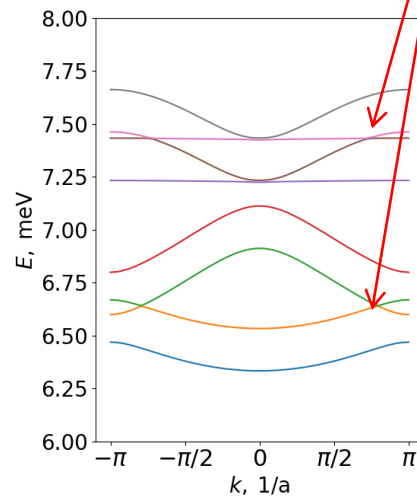
# 2D Arrays of polariton rings



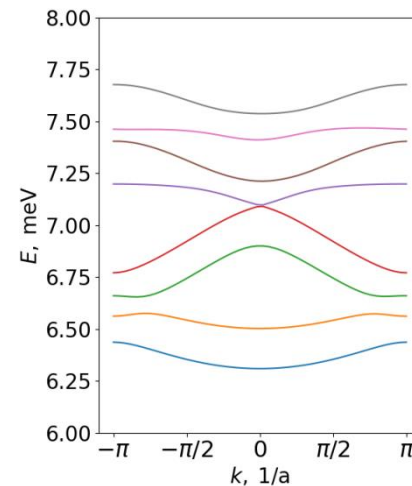
No magnetic fields



Zeeman only



Zeeman+ TE-TM



Bandgap openings:  
formation of  
topological edge  
modes

Thank you for attention

Спасибо за внимание

Ég þakka ykkur fyrir áheyrnina

